

Cohl Furey's $\mathbb{R} \otimes \mathbb{C} \otimes \mathbb{H} \otimes \mathbb{O}$ papers:

■ A Computational Analysis by J. Gregory Moxness

Motivation

This is an **updated** analysis of several papers from Cohl Furey related to using complexified real octonions (aka. bi-octonions) in the assignment of Standard Model (SM) Boson and Fermion particles to Clifford algebras and group symmetries. The previous analysis was based on [1-4]:

- [1] <https://arxiv.org/abs/1002.1497v5>
- [2] <https://arxiv.org/abs/1405.4601v1>
- [3] <https://arxiv.org/abs/1603.04078>
- [4] <https://arxiv.org/abs/1806.00612>
- [5] <https://arxiv.org/abs/1910.08395v1>

This .pdf is a direct output from my Mathematica (MTM) Notebook. I will follow up with a LaTex paper on the topic soon.

This notebook has code built in to operate symbolically on native MTM reals (\mathbb{R}), complexes ($\mathbb{C}=a+ib$), and quaternionic ($\mathbb{H}=a+ib+cj+dk$) forms, as well as my custom code to handle the octonions ($\mathbb{O}=\{e_0=1, e_1=i, e_2=j, e_3=k, e_4, e_5, e_6, e_7\}$), and now the bi-octonions (which doesn't assign the octonion e_1 to be equivalent to the complex imaginary (i)). That change also applies to the native quaternion assignments where of $e_1 = i$, $e_2 = j$, $e_3 = k$) in order to work with quater-octonions. This was a fairly trivial change to make since it simply involves removing the conversion of complex (and quaternion) operators from being involved in the octonionic multiplication (i.e. `SmallCircle[]` or symbolically \circ).

Commenting out that code is shown here:

Please note that my previous analysis here made the mistake of not commenting out these operations. As such, it was operating on octonions (not complexified bi-octonions), so some of my concerns were resolved based on correcting that error.

This also creates a split in the way conjugation is handled. The MTM native conjugation (`[Conjugate]` or `*`) handles the imaginary components of \mathbb{R} , \mathbb{C} , and \mathbb{H} (after loading `<<Quaternions``). With the change above, my code operates on octonion symbols (i.e. `oct = {e0, e1, e2, e3, e4, e5, e6, e7}`) as real as required for bi-octonions, but my octonion conjugation (`octConjugate[]` or $x^{(\text{asterisk})}$) computation operates on octonion symbols as imaginary in addition to the imaginary $\mathbb{C} \in \mathbb{H}$ elements if present (i.e. $i e_n \rightarrow (-i) (-e_n) = i e_n$).

The following code implements the two forms of conjugation introduced in the papers (`*` and `†`):

I've set up this "bi-octonion $\mathbb{R} \otimes \mathbb{C} \otimes \mathbb{H} \otimes \mathbb{O}$ machine" to easily switch from native MTM conjugation and my octonion conjugation.

The exceptions for `s` & `S` in this code accommodates their use in these papers as complex elements that only get MTM native treatment, thus avoiding recursive loops when not being reduced with used with the real octonion elements).

The following code implements the `s` and `S` in [5]:

```
In[2197]:= doConjugate = False;
```

```
In[2198]:= (* Conjugate *)
sScon = {
{s*, S},
{S, S*},
{s*, S*},
{s, S}};
```

In[2199]:=

It is also interesting to note that unlike complexes (\mathbb{C}), there are actually two forms of quaternion (\mathbb{H}) assignments where $k=(ij \text{ or } ji)$. For octonions, there are 480 possible multiplication tables that are non-trivial (i.e. not simple permutations of the oct symbols). There are 3840 split octonion multiplication tables. These multiplication tables are generated in my code using any of the algorithmically generated 480 possible octonion multiplication tables, as well as visualizing their Fano plane (or Fano cube) mnemonic based on which of the 7 triads selected. This code is also built to accommodate sedenions by Cayley-Dickson doubling, but I digress.

I am very interested here in the suggestion at the very end of Cohl's last paper [5] in the Addendum Section IX(B/C) on Multi-actions splitting spinor spaces, Lie algebras/groups, and Jordan algebras. I suspect having the ability to create a machine (i.e. a symbolic engine such as MTM) to operate on and visualize these structures as hyper-dimensional physical elements is critical to making progress in understanding our Universe more thoroughly.

From [1] Towards a unified theory of ideals (2010)

Select and implement the specific octonion:

So to set up the analysis, we select the Fano plane used in these papers, which happens to be the same one used by Baez in his work, which is in position 224 "flip \rightarrow True" in my assignment of the 480 combinations. The algorithm to visualize the Fano plane constructs it differently than the paper, but the triads are correct.

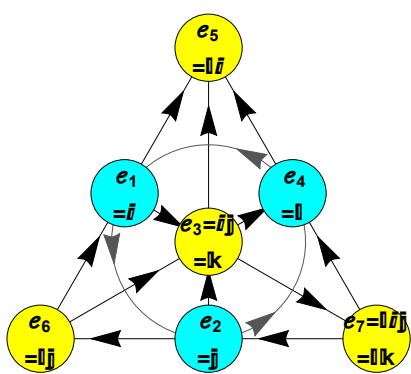
```
In[2204]:= (* Find Cohl Furey's Octonion = Baez's octonion based on triads *)
Position[(setFM[#, True];
StandardForm[triads]) & /@ Range@256,
{{1, 2, 4}, {1, 3, 7}, {1, 5, 6}, {2, 3, 5}, {2, 6, 7}, {3, 4, 6}, {4, 5, 7}}]
```

Out[2204]= (224 1)

```
In[2205]:= (* Set the Fano Multiplication matrix to that triad set (224, flip=True)
and display the triads, Fano plane and multiplication table *)
setFM[224, flip = True];
Row@{triads, fanoPlane,
 IJKLstyle = False; fmDispE,
 IJKLstyle = True;
 fmDispE}
```

$$\begin{pmatrix} 1 & 2 & 4 \\ 1 & 3 & 7 \\ 1 & 5 & 6 \\ 2 & 3 & 5 \\ 2 & 6 & 7 \\ 3 & 4 & 6 \\ 4 & 5 & 7 \end{pmatrix}$$

Octonion Fano Plane



$$\left(\begin{array}{ccccccc} e_0 \mapsto 1 & e_1 & e_2 & e_3 & e_4 & e_5 & e_6 & e_7 \\ e_1 & -1 & e_4 & e_7 & -e_2 & e_6 & -e_5 & -e_3 \\ e_2 & -e_4 & -1 & e_5 & e_1 & -e_3 & e_7 & -e_6 \\ e_3 & -e_7 & -e_5 & -1 & e_6 & e_2 & -e_4 & e_1 \\ e_4 & e_2 & -e_1 & -e_6 & -1 & e_7 & e_3 & -e_5 \\ e_5 & -e_6 & e_3 & -e_2 & -e_7 & -1 & e_1 & e_4 \\ e_6 & e_5 & -e_7 & e_4 & -e_3 & -e_1 & -1 & e_2 \\ e_7 & e_3 & e_6 & -e_1 & e_5 & -e_4 & -e_2 & -1 \end{array} \right) \left(\begin{array}{ccccccccc} e_0 \mapsto 1 & i & j & k & \mathbb{1} & \mathbb{i} & \mathbb{j} & \mathbb{k} \\ i & -1 & \mathbb{1} & \mathbb{k} & -\mathbb{j} & \mathbb{l}_j & -\mathbb{l}_i & -\mathbb{k} \\ j & -\mathbb{1} & -1 & \mathbb{i} & i & -\mathbb{k} & \mathbb{l}_k & -\mathbb{l}_j \\ k & -\mathbb{k} & -\mathbb{i} & -1 & \mathbb{l}_j & j & -\mathbb{l} & i \\ \mathbb{1} & \mathbb{j} & -i & -\mathbb{l}_j & -1 & \mathbb{l}_k & k & -\mathbb{l}_i \\ \mathbb{i} & -\mathbb{l}_j & k & -j & -\mathbb{l}_k & -1 & i & \mathbb{l} \\ \mathbb{j} & \mathbb{l}_i & -\mathbb{l}_k & \mathbb{l} & -\mathbb{k} & -i & -1 & j \\ \mathbb{k} & \mathbb{k} & \mathbb{l}_j & -i & \mathbb{l}_i & -\mathbb{l} & -j & -1 \end{array} \right)$$

It should be noted that there can be risk in making assumptions based on operating on a given set of triads → Fano planes → multiplication tables. Some of the patterns are specific only to that selected octonion out of the 480 (e.g. A reference in [2] from a reference to the Baez's paper (<https://arxiv.org/abs/math/0105155>) suggesting that index cycling and doubling patterns hold. While interesting, this is only true for that specific octonion multiplication table used in these papers.

Cohl's papers do not reduce the octonionic products (i.e. $e_1 \circ e_2 = e_{12} = e_4$) based on the selected multiplication table outside of using identities common to all imaginary numbers (e.g. $e_n \circ e_m = e_{nm}$, $e_{nm} = -e_{mn}$, and $e_n \circ e_n = -1$).

The following code implements functions needed for basic octonion processing (with some new code for bi-octonions):

Beginning the introduction of octonion chains including octonion reduction using the selected triad multiplication table.

```
In[2214]:= (* Some simple octProduct math *)
```

```
In[2215]:= e1 \circ e2
```

```
Out[2215]= e4
```

```
In[2216]:= e2 \circ e3
```

```
Out[2216]= e5
```

```
In[2217]:= e1 \circ (e2 \circ e3)
```

```
Out[2217]= e6
```

In[2218]:= $e_3 \circ (e_2 \circ e_1)$

Out[2218]= $-e_6$

In[2219]:= $e_6 + i \circ e_2$

Out[2219]= $e_4 + e_6$

Implementing Left vs. Right chains and precedence mapping function (f):

I constructed a function f to build the left (L) parenthetical multiplication precedence, which can be reversed by simply providing a “right→True” parameter to perform the right (R) operations. In addition, it takes a “right→both” parameter to both reverse the order and the parenthetical precedence. Note the use of the SmallCircle “◦” for doing proper octonion multiplication. While the paper does not show the evaluation of the octonionic multiplication, I do that in order to verify the logic.

```
In[2220]:= (* Define the parenthetical reverse nesting with Octonion multiplication table operation *)
Clear@f;
f[in_List, right_: False] := Module[{inFlat = Flatten@in, out = 1},
  Do[out = If[right == both, inFlat[[l]]◦out,
    If[right, out◦inFlat[[l]],
      Reverse[inFlat][[l]]◦out]], {l, Length@inFlat}];
  Simplify@out];
```

Using an example from [2] Generations: Three Prints, in Colour (2014) paragraph “Octonionic Chains”, we evaluate this function first w/o a defined octonion product table to see the proper parenthetical precedence is obtained for left and right operations:

In[2222]:= Clear[SmallCircle]

In[2223]:= f@{e3, e4, e6 + i◦e2}

Out[2223]= $e_3 \circ (e_4 \circ ((e_6 + i \circ e_2) \circ 1))$

Redefine SmallCircle and re-evaluate:

```
In[2224]:= SmallCircle[a_, 1] := a;
SmallCircle[1, b_] := b;
SmallCircle[a_List, b_List] := octonion@octProduct[a, b];
SmallCircle[a_, b_] := If[chksS@{a, b}, a b, oct2List@a◦oct2List@b];
```

```
In[2228]:= f@{e3, e4, e6 + i◦e2}
(* or in Mathematica postfix notation as a right to left functor as used in the paper *)
{e3, e4, e6 + i◦e2} // f
```

Out[2228]= $-1 + i e_7$

Out[2229]= $-1 + i e_7$

The associator function:

```
In[2230]:= associator[a_List, b_List, c_List] := (a◦b)◦c - a◦(b◦c);
associator[a_, b_, c_] := associator{oct2List@a, oct2List@b, oct2List@c};
```

If the associator is zero, it doesn’t show the non-associative property.

In[2232]:= associator[e3, e4, e5]

Out[2232]= $-2 e_1$

In[2233]:= associator[e3, e4, e6]

Out[2233]= 0

From [2]: Generations: Three Prints, in Colour (2014)

The following implements / checks the ν identities:

```
In[2234]:=  $\nu = -\frac{1}{2} f[\{e_3, e_4, e_6 + i e_2\}, \text{True}]$ 
Out[2234]=  $\frac{1}{2}(1 + i e_7)$ 

In[2235]:= (* Normal Conjugation - only negates  $i$  without changing  $e_n$  *)
 $\nu^*$ 
Out[2235]=  $\frac{1}{2}(1 - i e_7)$ 

In[2236]:= (* Cross Conjugation -
negates the real compared to normal (or both compared to octConjugation *)
 $\nu^\dagger$ 
Out[2236]=  $-\frac{1}{2} - \frac{i e_7}{2}$ 

In[2237]:= (* octConjugation - negates both  $i$  and  $e_n$  resulting in no change on  $i$   $e_n$  *)
 $\nu^*$ 
Out[2237]=  $\frac{1}{2} + \frac{i e_7}{2}$ 

In[2238]:= (* No Conjugation Identities *)
Simplify[ $\frac{1}{2}(1 + i e_7) == \nu == \nu \circ \nu$ ]
Out[2238]= True

In[2239]:= (* octConjugation Identities *)
Simplify[ $\frac{1}{2}(1 + i e_7) == \nu^* == \nu^* \circ \nu^*$ ]
Out[2239]= True

In[2240]:= (* Normal Conjugation Identities *)
Simplify[ $\frac{1}{2}(1 - i e_7) == \nu^* == \nu^* \circ \nu^*$ ]
 $0 == \nu \circ \nu^* == \nu^* \circ \nu$ 
Out[2240]= True

Out[2241]= True

In[2242]:= (* Cross Conjugation Identities *)
 $\nu^\dagger$ 
In[2243]:= 
Out[2243]=  $-\frac{1}{2} - \frac{i e_7}{2}$ 

In[2244]:=  $\nu^\dagger \circ \nu^\dagger$ 
Out[2244]=  $\frac{1}{2} + \frac{i e_7}{2}$ 
```

In[2245]:= (* Cross Conjugation Identities *)

$$\frac{1}{2} + \frac{\mathbf{i} e_7}{2} == -\mathbf{v}^\dagger == \mathbf{v}^\dagger \circ \mathbf{v}^\dagger$$

$$\mathbf{v}^\dagger == \mathbf{v} \circ \mathbf{v}^\dagger == \mathbf{v}^\dagger \circ \mathbf{v}$$

Out[2245]= True

Out[2246]= True

In[2247]:= $\mathbf{v} \odot \mathbf{v}^\dagger$

Out[2247]= 0

From [3]: Charge quantization from a number operator (2016)

The following implements / checks the new basis $\{\alpha_1, \alpha_2, \alpha_3\}$:

In[2248]:= Clear[\alpha, \alpha1, \alpha2, \alpha3]

In[2249]:= $\alpha_1 = \frac{1}{2} (-e_5 + \mathbf{i} e_4)$

Out[2249]= $\frac{1}{2} (-e_5 + i e_4)$

In[2250]:= (* Normal Conjugation *)

α_1^*

Out[2250]= $\frac{1}{2} (-e_5 - i e_4)$

In[2251]:= (* octConjugation *)

α_1^*

Out[2251]= $\frac{e_5}{2} + \frac{i e_4}{2}$

In[2252]:= (* Cross Conjugation *)

α_1^\dagger

Out[2252]= $-\frac{e_5}{2} - \frac{i e_4}{2}$

In[2253]:= $\alpha_2 = \frac{1}{2} (-e_3 + \mathbf{i} e_1)$

Out[2253]= $\frac{1}{2} (-e_3 + i e_1)$

In[2254]:= (* Normal Conjugation *)

α_2^*

Out[2254]= $\frac{1}{2} (-e_3 - i e_1)$

In[2255]:= α_2^*

Out[2255]= $\frac{e_3}{2} + \frac{i e_1}{2}$

In[2256]:= α_2^\dagger

Out[2256]= $-\frac{e_3}{2} - \frac{i e_1}{2}$

In[2257]:= $\alpha_3 = \frac{1}{2} (-e_6 + \mathbf{i} e_2)$

Out[2257]= $\frac{1}{2} (-e_6 + i e_2)$

In[2258]:= (* Normal Conjugation *)

α_3^*

$$\text{Out}[2258]= \frac{1}{2} (-e_6 - i e_2)$$

In[2259]:= α_3^*

$$\text{Out}[2259]= \frac{e_6}{2} + \frac{i e_2}{2}$$

In[2260]:= α_3^\dagger

$$\text{Out}[2260]= -\frac{e_6}{2} - \frac{i e_2}{2}$$

In[2261]:= Simplify[f@{\alpha1, \alpha2} == -I \alpha3^\dagger]

Simplify[f@{\alpha3, \alpha1} == -I \alpha2^\dagger]

Simplify[f@{\alpha2, \alpha3} == -I \alpha1^\dagger]

Out[2261]= True

Out[2262]= True

Out[2263]= True

Applying f and octonion multiplication reduces the following to zero:

In[2264]:= Simplify@{-{\alpha2^\dagger, \alpha1}, -{\alpha1^\dagger, \alpha2}}

{-f@{\alpha2^\dagger, \alpha1}, -f@{\alpha1^\dagger, \alpha2}}

$$\text{Out}[2264]= \begin{pmatrix} \frac{1}{2}(i e_1 + e_3) & \frac{1}{2}(e_5 - i e_4) \\ \frac{1}{2}(i e_4 + e_5) & \frac{1}{2}(e_3 - i e_1) \end{pmatrix}$$

$$\text{Out}[2265]= \{0, 0\}$$

In[2266]:= {#[[1]] \circ #[[2]]} & /@ {{\alpha1, \alpha2}, {\alpha2, \alpha3}, {\alpha1, \alpha3}}

{#[[1]]^\dagger \circ #[[2]]^\dagger} & /@ {{\alpha1, \alpha2}, {\alpha2, \alpha3}, {\alpha1, \alpha3}}

{#[[1]], #[[2]]^\dagger, #[[1]] \circ #[[2]]^\dagger, #[[2]] \circ #[[1]]^\dagger} & /@ {{\alpha1, \alpha2}, {\alpha2, \alpha3}, {\alpha1, \alpha3}}

$$\text{Out}[2266]= \begin{pmatrix} \frac{i e_6}{2} - \frac{e_2}{2} \\ \frac{i e_5}{2} - \frac{e_4}{2} \\ \frac{e_1}{2} - \frac{i e_3}{2} \end{pmatrix}$$

$$\text{Out}[2267]= \begin{pmatrix} -\frac{e_2}{2} - \frac{i e_6}{2} \\ -\frac{e_4}{2} - \frac{i e_5}{2} \\ \frac{e_1}{2} + \frac{i e_3}{2} \end{pmatrix}$$

$$\text{Out}[2268]= \begin{pmatrix} \frac{1}{2}(i e_4 - e_5) & -\frac{i e_1}{2} - \frac{e_3}{2} & 0 & 0 \\ \frac{1}{2}(i e_1 - e_3) & -\frac{i e_2}{2} - \frac{e_6}{2} & 0 & 0 \\ \frac{1}{2}(i e_4 - e_5) & -\frac{i e_2}{2} - \frac{e_6}{2} & 0 & 0 \end{pmatrix}$$

The following implements / checks ω identities:

In[2269]:= (* ω (or is it ν) is a negative charge particle that converts up \rightarrow down via right multiplication (i.e. W-) *)

$\omega = f@{\alpha1, \alpha2, \alpha3}$

$$\text{Out}[2269]= \frac{1}{2} (-e_7 + i)$$

In[2270]:= (* octConjugation *)

ω^*

$$\text{Out}[2270]= \frac{e_7}{2} - \frac{i}{2}$$

```
In[2271]:= (* Normal Conjugation *)
 $\omega^*$ 
Out[2271]=  $\frac{1}{2}(-e_7 - i)$ 

In[2272]:=  $\omega^\dagger$ 
Out[2272]=  $-\frac{e_7}{2} + \frac{i}{2}$ 

In[2273]:= (*  $\omega^\dagger$  (or is it  $\nu^*$ ) is a positive charge particle that converts down →
up via right multiplication (i.e.  $W+$ ) *)
Expand[ $\omega = \omega^\dagger$ ]

Out[2273]= True

In[2274]:= f@{ $\alpha_1^\dagger, \alpha_2^\dagger, \alpha_3^\dagger$ }
Out[2274]=  $\frac{1}{2}(-e_7 - i)$ 

In[2275]:=  $\nu = \text{Simplify}[-f@{\omega, \omega^\dagger}]$ 
Out[2275]=  $\frac{1}{2}(1 + i e_7)$ 

In[2276]:= Expand[- $\nu = f@{\omega, \omega^\dagger} = f@{\omega^\dagger, \omega} = \nu^\dagger$ ]
Out[2276]= True

In[2277]:= {ToString[Subscript["2\alpha", #], TraditionalForm] <>
    "=" <> ToString[Expand[2  $\alpha_\#$ ], TraditionalForm], Spacer@10,
    ToString[Subsuperscript["2\alpha", #, "t"], TraditionalForm] "=" <>
    ToString[Subsuperscript["2\alpha", #, "*"], TraditionalForm] <> "=" <>
    ToString[Subsuperscript["2\alpha", #, "oct*"], TraditionalForm] <>
    "=" <> ToString[Expand[2  $\alpha_\#^*$ ], TraditionalForm]} & /@ Range@3
Out[2277]=  $\begin{cases} 2\alpha_1 = 2\alpha_1 & 2\alpha_1^\dagger = 2\alpha_1^* = 2e_1 & 2\alpha_1^{\text{oct}*} = -2e_1 \\ 2\alpha_2 = 2\alpha_2 & 2\alpha_2^\dagger = 2\alpha_2^* = 2e_2 & 2\alpha_2^{\text{oct}*} = -2e_2 \\ 2\alpha_3 = 2\alpha_3 & 2\alpha_3^\dagger = 2\alpha_3^* = 2e_3 & 2\alpha_3^{\text{oct}*} = -2e_3 \end{cases}$ 

In[2278]:=  $\omega^\dagger \circ \alpha_\#^\dagger$  & /@ Range@3
Out[2278]=  $\left\{-\frac{e_3}{2} + \frac{ie_1}{2}, -\frac{e_6}{2} + \frac{ie_2}{2}, \frac{e_1}{2} + \frac{ie_3}{2}\right\}$ 

In[2279]:=  $\alpha_\# \circ \nu$  & /@ Range@3
Out[2279]=  $\left\{\frac{e_1}{2} - \frac{ie_3}{2}, \frac{e_2}{2} - \frac{ie_6}{2}, \frac{e_3}{2} + \frac{ie_1}{2}\right\}$ 

The following implements / checks the new basis  $\Lambda_1 - \Lambda_8$ :
Commutation (x○y) and AntiCommutation (xΘy)

In[2302]:= commutator[x_, y_] := Simplify@Chop[x○y - y○x] /. slRep
If[! doClifford,
  CircleDot[x_, y_] := commutator[x, y];
  derivation[x_, y_][a_] := (x○y)○a - 3 associator[x, y, a],
  derivation[x_, y_][a_] := commutator[commutator[x, y], a] - 3 associator[x, y, a]] /. slRep;

In[2304]:= antiCommutator[x_, y_] := Simplify@Chop[x○y + y○x] /. slRep;
CircleMinus[x_, y_] := antiCommutator[x, y];

Commutation checks
```

$$\text{In}[2306]:= \sum_{j=2}^6 \sum_{i=1}^{j-1} f @ \text{oct}[[\{i, j\}]] \odot f @ \text{oct}[[\{j, i\}]]$$

Out[2306]= 0

$$\text{In}[2307]:= \sum_{j=2}^6 \sum_{i=1}^{j-1} (f @ \text{oct}[[\{i, j\}]] + f @ \text{oct}[[\{j, i\}]])$$

Out[2307]= 0

AntiCommutation checks

$$\text{In}[2308]:= \alpha 1 \Theta \alpha 1$$

Out[2308]= 0

$$\text{In}[2309]:= \alpha 1 \Theta \alpha 2^\dagger$$

Out[2309]= 0

$$\text{In}[2310]:= \alpha 1 \Theta \alpha 1^\dagger$$

Out[2310]= -1

From [4]: $(SU(3))_C \times (SU(2))_L \times (U(1))_Y \times (U(1))_X$ as a symmetry of division algebraic ladder operators (2018)

4.3 Quarks and leptons as minimal left ideals (A one generation SM)

From [5] : Three generations, two unbroken gauge symmetries, and one eight-dimensional algebra (2019)

4.2 Clifford algebraic structure (Mf)

The following visualizes the 0-V to 6-V derived from MTM subsets:

```
In[2404]:= gather = GatherBy[Subsets@Range@6, Length@# &]
Length /@ %
```

```
Out[2405]= {1, 6, 15, 20, 15, 6, 1}
```

```
In[2406]:= invGather = (Complement[Range@7, #] & /@ #) & /@ gather
```

```

Out[2406]= {(1 2 3 4 5 6 7), {2 3 4 5 6 7}, {1 3 4 5 6 7}, {1 2 4 5 6 7}, {1 2 3 5 6 7}, {1 2 3 4 6 7}, {1 2 3 4 5 7}, {3 4 5 6 7}, {2 4 5 6 7}, {2 3 5 6 7}, {2 3 4 6 7}, {2 3 4 5 7}, {1 4 5 6 7}, {1 3 5 6 7}, {1 3 4 6 7}, {1 3 4 5 7}, {1 2 5 6 7}, {1 2 4 6 7}, {1 2 4 5 7}, {1 2 3 6 7}, {1 2 3 5 7}, {1 2 3 4 7}, {4 5 6 7}, {3 5 6 7}, {3 4 6 7}, {3 4 5 7}, {2 5 6 7}, {2 4 6 7}, {2 4 5 7}, {2 3 6 7}, {2 3 5 7}, {2 3 4 7}, {1 5 6 7}, {1 4 6 7}, {1 4 5 7}, {1 3 6 7}, {1 3 5 7}, {1 3 4 7}, {1 2 6 7}, {1 2 5 7}, {1 2 4 7}, {1 2 3 7}, {5 6 7}, {4 6 7}, {4 5 7}, {3 6 7}, {3 5 7}, {3 4 7}, {2 6 7}, {2 5 7}, {2 4 7}, {2 3 7}, {2 2 7}, {1 7} }, (7)

```

```
In[2407]:= (* Now show these as reduced bi-octonion elements L/R→Bottom/Top *)
addI = (*) False**True(**);
Table[{  

  If[addI && EvenQ@i, i, e0] o f@oct@gather[i, j]],  

  If[addI && EvenQ@i, i, e0] o f@oct@invGather[i, j]], {i, 7}, {j, Length@gather[i]}] /. e0 → 1
Reverse[Total@# & /@%]
(*)
  Total@#&/@%/ColumnForm
  Total@%/ColumnForm**)
```

$$\text{Out[2408]= } \left\{ (1 \quad -1), \begin{pmatrix} ie_1 & ie_1 \\ ie_2 & -ie_2 \\ ie_3 & ie_3 \\ ie_4 & -ie_4 \\ ie_5 & ie_5 \\ ie_6 & -ie_6 \end{pmatrix}, \begin{pmatrix} e_4 & e_4 \\ e_7 & -e_7 \\ -e_2 & -e_2 \\ e_6 & -e_6 \\ -e_5 & -e_5 \\ e_5 & e_5 \\ e_1 & -e_1 \\ e_7 & -e_7 \\ e_6 & e_6 \\ e_2 & -e_2 \\ -e_4 & -e_4 \\ e_7 & e_7 \\ e_3 & -e_3 \\ e_1 & e_1 \end{pmatrix}, \begin{pmatrix} ie_6 & -ie_6 \\ -i & -i \\ -ie_7 & ie_7 \\ -ie_3 & -ie_3 \\ -ie_5 & ie_5 \\ ie_4 & ie_4 \\ ie_2 & -ie_2 \\ -ie_3 & ie_3 \\ ie_7 & ie_7 \\ -i & i \\ ie_7 & ie_7 \\ -i & i \\ -ie_1 & -ie_1 \\ -ie_6 & -ie_6 \\ ie_5 & -ie_5 \\ -ie_4 & -ie_4 \\ ie_1 & -ie_1 \\ -i & -i \\ -ie_7 & ie_7 \\ ie_2 & ie_2 \end{pmatrix}, \begin{pmatrix} -e_3 & e_3 \\ -e_1 & -e_1 \\ 1 & -1 \\ e_5 & -e_5 \\ e_6 & e_6 \\ e_2 & -e_2 \\ -1 & -1 \\ -e_1 & e_1 \\ e_3 & e_3 \\ e_4 & -e_4 \\ -e_4 & e_4 \\ -e_2 & -e_2 \\ e_6 & -e_6 \\ -1 & -1 \\ -e_5 & e_5 \end{pmatrix}, \begin{pmatrix} ie_2 & ie_2 \\ -ie_4 & ie_4 \\ -ie_5 & -ie_5 \\ -ie_1 & ie_1 \\ -ie_6 & -ie_6 \\ ie_3 & -ie_3 \end{pmatrix}, (e_7 \quad e_7) \right\}$$

$$\text{Out[2409]= } \begin{pmatrix} e_7 & e_7 \\ -ie_1 + ie_2 + ie_3 - ie_4 - ie_5 - ie_6 & ie_1 + ie_2 - ie_3 + ie_4 - ie_5 - ie_6 \\ -2e_1 + 2e_6 - 1 & -2e_2 + 2e_3 - 3 \\ 2ie_2 - 2ie_3 - 4i & -2ie_1 - 2ie_6 + 4ie_7 \\ 2e_1 + 2e_6 + 3e_7 & -2e_2 - 2e_3 - e_7 \\ ie_1 + ie_2 + ie_3 + ie_4 + ie_5 + ie_6 & ie_1 - ie_2 + ie_3 - ie_4 + ie_5 - ie_6 \\ 1 & -1 \end{pmatrix}$$

Notice that the symmetry of the structure is more complicated than before applying the octonion Product.

Associativity of Composition of Maps

```
In[2410]:= Clear[F, G, H]
In[2413]:= F := f@{e2, e3, #} /. f → fff &;
G := f@{e3, #} /. f → fff &;
H := f@{e4, e6, e7, #} /. f → fff &;
In[2416]:= {F@#, G@#, H@#} & /@ {e0}
f[Flatten@%] /. f → fff
f[Flatten@%%, True] /. f → fff
f[Flatten@%%%, both] /. f → fff
```

$$\text{Out[2416]= } (e_5 \quad e_3 \quad -e_1)$$

$$\text{Out[2417]= } e_4$$

$$\text{Out[2418]= } -e_4$$

$$\text{Out[2419]= } -e_4$$

```
In[2420]:= (* Left to Right *)
  (F@e0) // G) // H
```

```
Out[2420]= -e4
```

```
In[2421]:= (* Right to Left *)
  F@ (G@ (H@e0))
```

```
Out[2421]= -e4
```

Redefine Λ & g as functions: Λ_n for $n = 1 - 8$ g_n for $n = 9 - 14$ (note: similar structure with $\frac{i}{2}$ factor)

```
In[2422]:= Clear[\Lambda, \Lambda1, \Lambda2, \Lambda3, \Lambda4, \Lambda5, \Lambda6, \Lambda7, \Lambda8, g]
```

```
In[2423]:= \Lambda[n_, ff_] := (***) Total[(***) f@{\#, ff} & /@ \Lambda_n(**)] (**);
  g[n_, ff_] := (***) Total[(***) f@{g_n, ff} & /@ g_n(**)] (**);
```

```
In[2425]:= (* \Lambda1 \rightarrow -e6/2 *)
```

```
In[2426]:= \Lambda1 = \frac{i}{2} \{ {e1, e5}, -{e3, e4} \};
```

```
\Lambda1@ff_ := \Lambda[1, ff];
```

```
\Lambda1@e0 /. f \rightarrow fff
```

```
Out[2428]= -\frac{e6}{2}
```

```
In[2429]:= (* \Lambda1 f \rightarrow v *)
```

```
In[2430]:= \Lambda[1, v] /. f \rightarrow fff
```

```
Out[2430]= \frac{1}{4} i (e2 + i e6)
```

```
In[2431]:= (* \Lambda2 \rightarrow \pm e2/2 \rightarrow 0 *)
```

```
In[2437]:= (* \Lambda3 \rightarrow -e7/2 *)
```

```
In[2443]:= (* \Lambda4 \rightarrow \pm e3/2 \rightarrow 0 *)
```

```
In[2449]:= (* \Lambda5 \rightarrow -e1/2 *)
```

```
In[2455]:= (* \Lambda6 \rightarrow \pm e5/2 \rightarrow 0 *)
```

```
In[2461]:= (* \Lambda7 \rightarrow \mp e4/2 \rightarrow 0 *)
```

```
In[2467]:= (* \Lambda8 \rightarrow -e7/2 *)
```

```
In[2473]:= (* g9 \rightarrow -e6/12^2 *)
```

```
In[2479]:= (* g10 \rightarrow +e2/12^2 *)
```

```
In[2485]:= (* g11 \rightarrow +e3/12^2 *)
```

```
In[2491]:= (* g12 \rightarrow +e1/12^2 *)
```

```
In[2497]:= (* g13 \rightarrow +e5/24 *)
```

```
In[2503]:= (* g14 \rightarrow +e4/12^2 *)
```

Three generation particle models

Redefine raising & lowering operators $\{\alpha_1, \alpha_2, \alpha_3\}$ as functions of $n = 1 - 3$ and f

Define SU(1) Q as a function of α_n $n = 1 - 3$ and f

```
In[2521]:= Clear@\n
```

```
In[2522]:= su1Q := \sum_{i=1}^3 \frac{1}{3} f @ {\alpha_i^\dagger, \alpha_i, \#} &;
```

```
In[2523]:= su1Q@eθ /. f → fff
%†
%%*
Out[2523]=  $\frac{1}{2} i (e_7 + i)$ 
Out[2524]=  $\frac{1}{2} - \frac{i e_7}{2}$ 
Out[2525]=  $-\frac{1}{2} + \frac{i e_7}{2}$ 

In[2526]:= su1Q@e4
Simplify@%
Out[2526]=  $\frac{1}{3} f\left(\left\{-\frac{e_3}{2} - \frac{i e_1}{2}, \frac{1}{2}(-e_3 + i e_1), e_4\right\}\right) + \frac{1}{3} f\left(\left\{-\frac{e_5}{2} - \frac{i e_4}{2}, \frac{1}{2}(-e_5 + i e_4), e_4\right\}\right) + \frac{1}{3} f\left(\left\{-\frac{e_6}{2} - \frac{i e_2}{2}, \frac{1}{2}(-e_6 + i e_2), e_4\right\}\right)$ 
Out[2527]=  $\frac{1}{3} \left(f\left(\left\{-\frac{1}{2} i (e_1 - i e_3), \frac{1}{2} i (e_1 + i e_3), e_4\right\}\right) + f\left(\left\{-\frac{1}{2} i (e_4 - i e_5), \frac{1}{2} i (e_4 + i e_5), e_4\right\}\right) + f\left(\left\{-\frac{1}{2} i (e_2 - i e_6), \frac{1}{2} i (e_2 + i e_6), e_4\right\}\right)\right)$ 

In[2528]:= f@{α#†, α#} /. f → fff & /@ Range@3
Out[2528]=  $\left\{\frac{1}{2} i (e_7 + i), \frac{1}{2} i (e_7 + i), \frac{1}{2} i (e_7 + i)\right\}$ 

Define exp G as a function of j = 1 – 8 → rj and f

In[2529]:= Clear@G

In[2530]:= rθ = 2.;
r1 = r2 = r3 = r4 = r5 = r6 = r7 = r8 = 1.;
G[j_, ff_] := ei rj Δ[j, ff] + i rθ su1Q@ff ;

In[2533]:= i r1 Δ[1, 1] /. f → fff
i rθ su1Q@eθ /. f → fff
Out[2533]=  $(0. - 0.5 i) e_6$ 
Out[2534]=  $(-1. + 0. i) (e_7 + i)$ 

In[2535]:= G[1, 1] /. f → fff
Out[2535]=  $\exp((0. - 0.5 i) e_6 - (1. + 0. i) (e_7 + i))$ 

In[2536]:= Chop@{#, G[1, 1] α# G[1, 1]-1, G[1, 1] α#† G[1, 1]-1} & /@ Range@3 /. f → fff
Out[2536]=  $\begin{pmatrix} 1 & 0.5(i e_4 - e_5) & 1. \left(-\frac{i e_4}{2} - \frac{e_5}{2}\right) \\ 2 & 0.5(i e_1 - e_3) & 1. \left(-\frac{i e_1}{2} - \frac{e_3}{2}\right) \\ 3 & 0.5(i e_2 - e_6) & 1. \left(-\frac{i e_2}{2} - \frac{e_6}{2}\right) \end{pmatrix}$ 
```

Define the 3 generations of quarks and leptons (2014 and 2019 versions)

```
In[2537]:= Clear[u1, u2, u3, Λ1Λ2, Λ4Λ5, Λ6Λ7,
u1r, u1g, u1b, u2r, u2g, u2b, u3r, u3g, u3b,
d123rgb, d1r, d1g, d1b, d2r, d2g, d2b, d3r, d3g, d3b,
v1, v2, v3, e123, e1, e2, e3]

In[2538]:= u1@oct_ := {ss[[2]], Flatten@{Total@oct}, ss[[4]]};
u2@oct_ := {ss[[3]], Flatten@{Total@oct}, ss[[1]]};
u3@oct_ := {ss[[4]], Flatten@{Total@oct}, ss[[1]]};

In[2541]:= Δnm[nm_List] := f@# & /@ Flatten[{nm}, 1];

In[2542]:= (* u1r construction / deconstruction *)
```

In[2543]:= $-\text{ix} \Delta_6 \{-1, 1\} // \text{StandardForm}$

$\Delta_{\text{nm}}[-\text{ix} \Delta_6] \{-1, 1\}$

$\text{Total}@% /. f \rightarrow \text{fff}$

Out[2543]/StandardForm=

$$\left\{ \left\{ -\frac{e_1}{2}, -\frac{e_6}{2} \right\}, \left\{ \frac{e_2}{2}, \frac{e_3}{2} \right\} \right\}$$

Out[2544]= $\left\{ -f\left(\left\{ \frac{e_1}{2}, \frac{e_6}{2} \right\}\right), f\left(\left\{ \frac{e_2}{2}, \frac{e_3}{2} \right\}\right) \right\}$

$$\text{Out}[2545]= \frac{e_5}{2}$$

In[2546]:= $\Delta_7 \{-1, 1\} // \text{StandardForm}$

$\Delta_{\text{nm}}[\Delta_7] \{-1, 1\}$

$\text{Total}@% /. f \rightarrow \text{fff}$

Out[2546]/StandardForm=

$$\left\{ \left\{ -\frac{\text{ix} e_1}{2}, -\frac{\text{ix} e_2}{2} \right\}, \left\{ \frac{\text{ix} e_3}{2}, \frac{\text{ix} e_6}{2} \right\} \right\}$$

Out[2547]= $\left\{ -f\left(\left\{ \frac{ie_1}{2}, \frac{ie_2}{2} \right\}\right), f\left(\left\{ \frac{ie_3}{2}, \frac{ie_6}{2} \right\}\right) \right\}$

$$\text{Out}[2548]= \frac{e_4}{2}$$

In[2549]:= $\Delta_6 \Delta_7 = (\Delta_{\text{nm}}[-\text{ix} \Delta_6] + \Delta_{\text{nm}}[\Delta_7]) \{-1, 1\}$

% /. f → fff

$\text{Total}@%$

Out[2549]= $\left\{ -f\left(\left\{ \frac{e_1}{2}, \frac{e_6}{2} \right\}\right) - f\left(\left\{ \frac{ie_1}{2}, \frac{ie_2}{2} \right\}\right), f\left(\left\{ \frac{e_2}{2}, \frac{e_3}{2} \right\}\right) + f\left(\left\{ \frac{ie_3}{2}, \frac{ie_6}{2} \right\}\right) \right\}$

$$\text{Out}[2550]= \left\{ \frac{e_4}{4} + \frac{e_5}{4}, \frac{e_4}{4} + \frac{e_5}{4} \right\}$$

$$\text{Out}[2551]= \frac{e_4}{2} + \frac{e_5}{2}$$

Use this change for 2014 u1 R→G fix:

In[2552]:= $\Delta_6 \Delta_7 = (**)-2(*) \Delta_6 \Delta_7 = (\Delta_{\text{nm}}[\Delta_6] + \Delta_{\text{nm}}[\Delta_7]) \{-1, 1\}$

$\text{Expand}@% /. f \rightarrow \text{fff}$

$\text{Total}@%$

Out[2552]= $\left\{ 2\left(f\left(\left\{ \frac{ie_1}{2}, \frac{ie_2}{2} \right\}\right) + if\left(\left\{ \frac{ie_1}{2}, \frac{ie_6}{2} \right\}\right)\right), -2\left(if\left(\left\{ \frac{ie_2}{2}, \frac{ie_3}{2} \right\}\right) + f\left(\left\{ \frac{ie_3}{2}, \frac{ie_6}{2} \right\}\right)\right) \right\}$

$$\text{Out}[2553]= \left\{ -\frac{e_4}{2} + \frac{ie_5}{2}, -\frac{e_4}{2} + \frac{ie_5}{2} \right\}$$

$$\text{Out}[2554]= -e_4 + ie_5$$

In[2555]:= (* u1g construction / deconstruction *)

In[2556]:= $\text{ix} \Delta_4 \{1, -1\} // \text{StandardForm}$

$\Delta_{\text{nm}}[\text{ix} \Delta_4] \{1, -1\}$

$\text{Total}@% /. f \rightarrow \text{fff}$

Out[2556]/StandardForm=

$$\left\{ \left\{ -\frac{e_2}{2}, -\frac{e_5}{2} \right\}, \left\{ \frac{e_4}{2}, \frac{e_6}{2} \right\} \right\}$$

Out[2557]= $\left\{ f\left(\left\{ -\frac{e_2}{2}, -\frac{e_5}{2} \right\}\right), -f\left(\left\{ -\frac{e_4}{2}, -\frac{e_6}{2} \right\}\right) \right\}$

$$\text{Out}[2558]= -\frac{e_3}{2}$$

```
In[2559]:=  $\Delta_5 \{1, -1\} // \text{StandardForm}$ 
 $\Delta_{\text{nm}}[\Delta_5] \{1, -1\}$ 
 $\text{Total}@% /. f \rightarrow \text{fff}$ 
Out[2559]/StandardForm=
 $\left\{ \left\{ -\frac{\frac{1}{2}e_5}{2}, -\frac{\frac{1}{2}e_6}{2} \right\}, \left\{ -\frac{\frac{1}{2}e_2}{2}, -\frac{\frac{1}{2}e_4}{2} \right\} \right\}$ 
```

```
Out[2560]=  $\left\{ f\left(\left\{ -\frac{i e_5}{2}, -\frac{i e_6}{2} \right\}\right), -f\left(\left\{ \frac{i e_2}{2}, \frac{i e_4}{2} \right\}\right) \right\}$ 
```

```
Out[2561]= 0
```

```
In[2562]:=  $\Delta_4 \Delta_5 = (\Delta_{\text{nm}}[\dot{\Delta}_4] + \Delta_{\text{nm}}[\Delta_5]) \{1, -1\}$ 
% /. f \rightarrow \text{fff}
 $\text{Total}@%$ 
```

```
Out[2562]=  $\left\{ f\left(\left\{ -\frac{e_2}{2}, -\frac{e_5}{2} \right\}\right) + f\left(\left\{ -\frac{i e_5}{2}, -\frac{i e_6}{2} \right\}\right), -f\left(\left\{ -\frac{e_4}{2}, -\frac{e_6}{2} \right\}\right) - f\left(\left\{ \frac{i e_2}{2}, \frac{i e_4}{2} \right\}\right) \right\}$ 
```

```
Out[2563]=  $\left\{ -\frac{e_1}{4} - \frac{e_3}{4}, \frac{e_1}{4} - \frac{e_3}{4} \right\}$ 
```

```
Out[2564]=  $-\frac{e_3}{2}$ 
```

Use this change for 2014 u1 G→R fix:

```
In[2565]:=  $\Delta_4 \Delta_5 = (***) 2 (* 4 **) (\Delta_{\text{nm}}[\Delta_4] \{1, -1\} + \dot{\Delta}_{\text{nm}}[\Delta_5])$ 
 $\text{Expand}[% /. f \rightarrow \text{fff}]$ 
 $\text{Total}@%$ 
```

```
Out[2565]=  $\left\{ 2\left(f\left(\left\{ \frac{i e_2}{2}, \frac{i e_5}{2} \right\}\right) + i f\left(\left\{ -\frac{i e_5}{2}, -\frac{i e_6}{2} \right\}\right)\right), 2\left(i f\left(\left\{ \frac{i e_2}{2}, \frac{i e_4}{2} \right\}\right) - f\left(\left\{ \frac{i e_4}{2}, \frac{i e_6}{2} \right\}\right)\right) \right\}$ 
```

```
Out[2566]=  $\left\{ \frac{e_3}{2} - \frac{i e_1}{2}, \frac{e_3}{2} - \frac{i e_1}{2} \right\}$ 
```

```
Out[2567]=  $e_3 - i e_1$ 
```

```
In[2568]:= (* u1b construction / deconstruction *)
```

```
In[2569]:=  $\dot{\Delta}_1 \{-1, 1\} // \text{StandardForm}$ 
 $\Delta_{\text{nm}}[\dot{\Delta}_1] \{-1, 1\}$ 
 $\text{Total}@% /. f \rightarrow \text{fff}$ 
```

```
Out[2569]/StandardForm=
 $\left\{ \left\{ \frac{e_1}{2}, \frac{e_5}{2} \right\}, \left\{ \frac{e_3}{2}, \frac{e_4}{2} \right\} \right\}$ 
```

```
Out[2570]=  $\left\{ -f\left(\left\{ -\frac{e_1}{2}, -\frac{e_5}{2} \right\}\right), f\left(\left\{ \frac{e_3}{2}, \frac{e_4}{2} \right\}\right) \right\}$ 
```

```
Out[2571]= 0
```

```
In[2572]:=  $\Delta_2 \{-1, 1\} // \text{StandardForm}$ 
 $\Delta_{\text{nm}}[\Delta_2] \{-1, 1\}$ 
 $\text{Total}@% /. f \rightarrow \text{fff}$ 
```

```
Out[2572]/StandardForm=
 $\left\{ \left\{ \frac{\frac{1}{2}e_1}{2}, \frac{\frac{1}{2}e_4}{2} \right\}, \left\{ -\frac{\frac{1}{2}e_3}{2}, -\frac{\frac{1}{2}e_5}{2} \right\} \right\}$ 
```

```
Out[2573]=  $\left\{ -f\left(\left\{ -\frac{i e_1}{2}, -\frac{i e_4}{2} \right\}\right), f\left(\left\{ -\frac{i e_3}{2}, -\frac{i e_5}{2} \right\}\right) \right\}$ 
```

```
Out[2574]=  $-\frac{e_2}{2}$ 
```

```

In[2575]:= A1A2 =  $(\Delta nm[\dot{\Delta} \Delta_1] + \Delta nm[\Delta_2]) \{-1, 1\}$ 
% /. f → fff
Total@%

Out[2575]=  $\left\{ -f\left(\left\{-\frac{e_1}{2}, -\frac{e_5}{2}\right\}\right) - f\left(\left\{-\frac{ie_1}{2}, -\frac{ie_4}{2}\right\}\right), f\left(\left\{\frac{e_3}{2}, \frac{e_4}{2}\right\}\right) + f\left(\left\{-\frac{ie_3}{2}, -\frac{ie_5}{2}\right\}\right) \right\}$ 
Out[2576]=  $\left\{ -\frac{e_2}{4} - \frac{e_6}{4}, \frac{e_6}{4} - \frac{e_2}{4} \right\}$ 
Out[2577]=  $-\frac{e_2}{2}$ 

In[2578]:= (* u1 RGB *)

In[2579]:= u1r := u1[ $\Delta 6 \Delta 7$ ];
u1r
u1r /. f → fff

Out[2580]=  $\left\{ \{s, S^*\}, \left\{ 2\left(f\left(\left\{\frac{ie_1}{2}, \frac{ie_2}{2}\right\}\right) + if\left(\left\{\frac{ie_1}{2}, \frac{ie_6}{2}\right\}\right)\right) - 2\left(if\left(\left\{\frac{ie_2}{2}, \frac{ie_3}{2}\right\}\right) + f\left(\left\{\frac{ie_3}{2}, \frac{ie_6}{2}\right\}\right)\right) \right\}, \{s, S\} \right\}$ 
Out[2581]=  $\left\{ \{s, S^*\}, \left\{ 2\left(-\frac{e_4}{4} + \frac{ie_5}{4}\right) - 2\left(\frac{e_4}{4} - \frac{ie_5}{4}\right) \right\}, \{s, S\} \right\}$ 

In[2582]:= (* From 2014 paper *)
%[[2, 1]] ∘ ν

Out[2582]=  $-e_4 + ie_5$ 

In[2583]:= u1g := u1[ $\Delta 4 \Delta 5$ ];
u1g
u1g /. f → fff

Out[2584]=  $\left\{ \{s, S^*\}, \left\{ 2\left(if\left(\left\{\frac{ie_2}{2}, \frac{ie_4}{2}\right\}\right) - f\left(\left\{\frac{ie_4}{2}, \frac{ie_6}{2}\right\}\right)\right) + 2\left(f\left(\left\{\frac{ie_2}{2}, \frac{ie_5}{2}\right\}\right) + if\left(\left\{-\frac{ie_5}{2}, -\frac{ie_6}{2}\right\}\right)\right) \right\}, \{s, S\} \right\}$ 
Out[2585]=  $\left\{ \{s, S^*\}, \left\{ 4\left(\frac{e_3}{4} - \frac{ie_1}{4}\right) \right\}, \{s, S\} \right\}$ 

In[2586]:= (* From 2014 paper *)
%[[2, 1]] ∘ ν

Out[2586]= 0

In[2587]:= u1b := u1[ $\Delta 1 \Delta 2$ ];
u1b
u1b /. f → fff

Out[2588]=  $\left\{ \{s, S^*\}, \left\{ -f\left(\left\{-\frac{e_1}{2}, -\frac{e_5}{2}\right\}\right) - f\left(\left\{-\frac{ie_1}{2}, -\frac{ie_4}{2}\right\}\right) + f\left(\left\{-\frac{ie_3}{2}, -\frac{ie_5}{2}\right\}\right) + f\left(\left\{\frac{e_3}{2}, \frac{e_4}{2}\right\}\right) \right\}, \{s, S\} \right\}$ 
Out[2589]=  $\left\{ \{s, S^*\}, \left\{ -\frac{e_2}{2} \right\}, \{s, S\} \right\}$ 

In[2590]:= (* From 2014 paper *)
%[[2, 1]] ∘ ν

Out[2590]=  $-\frac{e_2}{4} + \frac{ie_6}{4}$ 

In[2591]:= (* u2 RGB *)

In[2604]:= (* u3 RGB *)

In[2624]:= (* d1 RGB *)

In[2625]:= d123rgb@oct_ := (* Row@**) {ss[[2]], Total@oct, ss[[1]]};

```

```

In[2626]:= d1r := d123rgb[{{f@{e1, e2, e6} - f@{e3}}, -{f@{e1, e4, e5} + i f@{e1}}}] ;
d1r
d1r /. f → fff

Out[2627]= {{s, S*}, {f({e1, e2, e6}) - f({e1, e4, e5}) - i f({e1}) - f({e3})}, {s*, S} }

Out[2628]= {{s, S*}, {-e3 - i e1}, {s*, S} }

In[2629]:= (* From 2014 paper *)
%[[2, 1]] ° ν

Out[2629]= -e3 - i e1

In[2630]:= d1g := d123rgb[{{f@{e1, e3, e4} + f@{e5}}, i {-f@{e1, e3, e5} + f@{e4}}}] ;
(* ~u3r *)
d1g
d1g /. f → fff

Out[2631]= {{s, S*}, {f({e1, e3, e4}) + i (f({e4}) - f({e1, e3, e5})) + f({e5})}, {s*, S} }

Out[2632]= {{s, S*}, {0}, {s*, S} }

In[2633]:= (* From 2014 paper *)
%[[2, 1]] ° ν

Out[2633]= 0

In[2634]:= d1b := d123rgb[{-{f@{e1, e2, e5} + f@{e1, e4, e6}}, i {f@{e1, e5, e6} - f@{e1, e2, e4}}}] ;
d1b
d1b /. f → fff

Out[2635]= {{s, S*}, {-f({e1, e2, e5}) - f({e1, e4, e6}) + i (f({e1, e5, e6}) - f({e1, e2, e4}))}, {s*, S} }

Out[2636]= {{s, S*}, {0}, {s*, S} }

In[2637]:= (* From 2014 paper *)
%[[2, 1]] ° ν

Out[2637]= 0

In[2638]:= (* d2 RGB *)
In[2651]:= (* d3 RGB *)
In[2664]:= (* ν *)
In[2665]:= ν1 := (* Row@**) {ss[[1]], Total@{1}, i {f@{e1, e3} + f@{e2, e6} + f@{e4, e5}}}, ss[[1]];
ν1 /. f → fff

Out[2666]= {{s*, S}, {1 + 3 i e7}, {s*, S} }

In[2667]:= (* From 2014 paper *)
ν1 := (* Row@**)
{ss[[1]], 1 + i 0.}, ss[[1]];
ν1 /. f → fff
%[[2, 1]] ° ν

Out[2668]= {{s*, S}, {1. + 0. i}, {s*, S} }

Out[2669]= 0.5 + (0. + 0.5 i) e7

In[2670]:= (*
    ν1 *)
ν2 := (* Row@**) {ss[[2]], Total@{3}, -i {f@{e1, e3} + f@{e2, e6} + f@{e4, e5}}}, ss[[2]];
ν2 /. f → fff

Out[2671]= {{s, S*}, {3 - 3 i e7}, {s, S*}}

```

```
In[2672]:= (* From 2014 paper *)
(*
v1 *)
v2 := (* Row@**) {ss[[2]], Total@{{f@{e1, e3} + f@{e2, e6} + f@{e4, e5}}}, ss[[2]]};
v2 /. f → fff
%[[2, 1]] ∘ v
```

Out[2673]= $\{\{s, S^*\}, \{3 e_7\}, \{s, S^*\}\}$

$$\frac{3 e_7}{2} - \frac{3 i}{2}$$

```
In[2675]:= (*
d1b *)
v3 := (* Row@**) {ss[[3]],
Total@{-{f@{e1, e2, e5} + f@{e1, e4, e6}}, -i {f@{e1, e5, e6} + f@{e1, e2, e4}}}, ss[[4]]};
v3 /.
f →
fff
```

Out[2676]= $\{\{s^*, S^*\}, \{2 i\}, \{s, S\}\}$

```
In[2677]:= (* From 2014 paper *)
%[[2, 1]] ∘ v
```

$$-e_7 + i$$

```
In[2678]:= (* v *)
```

Define SU(3) Q to match the action in eqs. 19, 22 and 23

While I've tried many variations, this one is simply based on the position in the list of sS conjugate combinations (given the leading sS pair in the defined particle).

```
In[2695]:= su3Q@n_ := Module[{in = n[[2]],
pos = position[sS, n[[1]]]},
(**) Simplify@(**) If[pos === {}, in, {(**) i, (**)
pos
3, n[[1, 1]]}(*)./
{S → 1/2 (f@{in} + 1/2 (-f@{e7, in} + f@{e1, e3, in} + f@{e2, e6, in} + f@{e4, e5, in})),
S → 1/2 (f@{in} + 1/2 (-f[{e7, in}, True] +
f[{e1, e3, in}, True] + f[{e2, e6, in}, True] + f[{e4, e5, in}, True]))}**)]];
```

This works for the example given in u1r,

but I can't figure out how to make it work across all flavors and generations given their particle definitions above. My confusion seems to be around the proper use of the sS components.

Do an example, showing the decomposition of sS and show su3Q taking u1r as input (with and without evaluating f):

```
In[2696]:= doConjugate = True;
```

```
In[2697]:= sS /. {s → 1/2 (f@{in} + 1/2 (-f@{e7, in} + f@{e1, e3, in} + f@{e2, e6, in} + f@{e4, e5, in})),
s → 1/2 (f@{in} + 1/2 (-f[{e7, in}, True] + f[{e1, e3, in}, True] +
f[{e2, e6, in}, True] + f[{e4, e5, in}, True])) /. {in → 1, f → fff}}
```

$$\begin{pmatrix} \frac{1}{2}(1 - ie_7) & \frac{1}{2}(ie_7 + 1) \\ \frac{1}{2}(ie_7 + 1) & \frac{1}{2}(1 - ie_7) \\ \frac{1}{2}(1 - ie_7) & \frac{1}{2}(1 - ie_7) \\ \frac{1}{2}(ie_7 + 1) & \frac{1}{2}(ie_7 + 1) \end{pmatrix}$$

In[2698]:= Row@u1r

```
Out[2698]= {s, S*} \{2 \left(f\left(\left\{\frac{i e_1}{2}, \frac{i e_2}{2}\right\}\right)+i f\left(\left\{\frac{i e_1}{2}, \frac{i e_6}{2}\right\}\right)\right)-2 \left(i f\left(\left\{\frac{i e_2}{2}, \frac{i e_3}{2}\right\}\right)+f\left(\left\{\frac{i e_3}{2}, \frac{i e_6}{2}\right\}\right)\right)\} \{s, S\}
```

```
In[2699]:= u1r[[2, 1]] /. f → fff
```

$$\text{Out}[2699] = 2 \left(-\frac{e_4}{4} + \frac{i e_5}{4} \right) - 2 \left(\frac{e_4}{4} - \frac{i e_5}{4} \right)$$

```
In[2700]:= Expand@#&[3]{f@{sS[[#,1]],f@{sS[[#,2]]}}}&.
```

$$\left\{ s \rightarrow \frac{1}{2} \left(f @ \{in\} + \frac{\frac{1}{2}}{2} (-f @ \{e_7, in\} + f @ \{e_1, e_3, in\} + f @ \{e_2, e_6, in\} + f @ \{e_4, e_5, in\}) \right) \right\},$$

$$S \rightarrow \frac{1}{2} \left(f @ \{in\} + \frac{\mathbf{i}}{2} (-f[\{e_7, in\}, True] + f[\{e_1, e_3, in\}, True]) + \right.$$

$$f[\{e_2, e_6, \text{in}\}, \text{True}] + f[\{e_4, e_5, \text{in}\}, \text{True}]) \Big) \Big\} /.$$

```
{in → u1r[2, 1] /. f → fff} & /@ Range@3
```

% /. f → fff

Expand@Total [Total /@ %]

Out[2700]=
$$\begin{cases} \frac{1}{3} f\left(\frac{1}{2}\left(f\left(2\left(\frac{i e_5}{4}-\frac{e_4}{4}\right)-2\left(\frac{e_4}{4}-\frac{i e_5}{4}\right)\right)\right)+\frac{1}{2} i\left(-f\left(\left\{e_7,2\left(\frac{i e_5}{4}-\frac{e_4}{4}\right)-2\left(\frac{e_4}{4}-\frac{i e_5}{4}\right)\right)\right)\right)+f\left(\left\{e_1,e_3,2\left(\frac{i e_5}{4}-\frac{e_4}{4}\right)-2\left(\frac{e_4}{4}-\frac{i e_5}{4}\right)\right)\right)+f\left(\left\{e_2,e_6,2\left(\frac{i e_5}{4}-\frac{e_4}{4}\right)-2\left(\frac{e_4}{4}-\frac{i e_5}{4}\right)\right)\right), \\ \frac{2}{3} f\left(\left\{\frac{1}{2}\left(f\left(2\left(\frac{i e_5}{4}-\frac{e_4}{4}\right)-2\left(\frac{e_4}{4}-\frac{i e_5}{4}\right)\right)\right)+\frac{1}{2} i\left(-f\left(\left\{e_7,2\left(\frac{i e_5}{4}-\frac{e_4}{4}\right)-2\left(\frac{e_4}{4}-\frac{i e_5}{4}\right)\right)\right)\right)+f\left(\left\{e_1,e_3,2\left(\frac{i e_5}{4}-\frac{e_4}{4}\right)-2\left(\frac{e_4}{4}-\frac{i e_5}{4}\right)\right)\right)+f\left(\left\{e_2,e_6,2\left(\frac{i e_5}{4}-\frac{e_4}{4}\right)-2\left(\frac{e_4}{4}-\frac{i e_5}{4}\right)\right)\right), \\ f\left(\left\{\frac{1}{2}\left(f\left(2\left(\frac{i e_5}{4}-\frac{e_4}{4}\right)-2\left(\frac{e_4}{4}-\frac{i e_5}{4}\right)\right)\right)+\frac{1}{2} i\left(-f\left(\left\{e_7,2\left(\frac{i e_5}{4}-\frac{e_4}{4}\right)-2\left(\frac{e_4}{4}-\frac{i e_5}{4}\right)\right)\right)\right)+f\left(\left\{e_1,e_3,2\left(\frac{i e_5}{4}-\frac{e_4}{4}\right)-2\left(\frac{e_4}{4}-\frac{i e_5}{4}\right)\right)\right)+f\left(\left\{e_2,e_6,2\left(\frac{i e_5}{4}-\frac{e_4}{4}\right)-2\left(\frac{e_4}{4}-\frac{i e_5}{4}\right)\right)\right) \end{cases}$$

$$\text{Out}[2701] = \begin{pmatrix} \frac{1}{3}(-e_4 - ie_5) & 0 \\ \frac{2}{3}(ie_5 - e_4) & 0 \\ -e_4 - ie_5 & 0 \end{pmatrix}$$

Out[2702]= $-2e_4 - \frac{2ie_5}{3}$

```
In[2703]:= doConjugate = False;
```

su3Q@u1r[2, 1]

$$\text{Out}[2704] = -2 \left(i f\left(\left\{\frac{i e_2}{2}, \frac{i e_3}{2}\right\}\right) + f\left(\left\{\frac{i e_3}{2}, \frac{i e_6}{2}\right\}\right) \right)$$

```
In[2705]:= doConjugate = True;
```

su3Q@u1r[2, 1]

$$\text{Out}[2706] = -2 \left(i f\left(\left\{\frac{i e_2}{2}, \frac{i e_3}{2}\right\}\right) + f\left(\left\{\frac{i e_3}{2}, \frac{i e_6}{2}\right\}\right) \right)$$

From the 2014 particle model in [2] show commutation action for U1R \rightarrow U1G with octonion multiplication

Cohl doesn't explicitly reduce the octonion elements via octonionic multiplication here, but does reduce them via common identities. Here we process the manual reduction of the color changing action of $\Delta_1[\nu] \odot u1r = \nu \circ u1g$ shown in eqs. 20 & 21 of [2], and then show the commutation directly.

So here we replicate the identity reductions, but also process each step with octonionic multiplication without the \circ .
 v :

```
In[2707]:= { -i f @ {e1, e5, e1, e2}, -f @ {e1, e5, e1, e6}, +f @ {e1, e5, e2, e3}, +i f @ {e1, e5, e3, e6} } /. f → ffff
```

Out[2707]= $\{-i e_3, -e_1, -e_1, -i e_3\}$

Out[2708]= $-2 e_1 - 2 i e_3$

```
In[2709]:= {+i f@{e3, e4, e1, e2}, +f@{e3, e4, e1, e6}, -f@{e3, e4, e2, e3}, -i f@{e3, e4, e3, e6}} /. f → fff
Total@%
```

```
Out[2709]= {-i e3, -e1, -e1, -i e3}
```

```
Out[2710]= -2 e1 - 2 i e3
```

```
In[2711]:= {+i f@{e1, e2, e1, e5}, +f@{e1, e6, e1, e5}, -f@{e2, e3, e1, e5}, -i f@{e3, e6, e1, e5}} /. f → fff
Total@%
```

```
Out[2711]= {-i e3, -e1, e1, i e3}
```

```
Out[2712]= 0
```

```
In[2713]:= {-i f@{e1, e2, e3, e4}, -f@{e1, e6, e3, e4}, +f@{e2, e3, e3, e4}, +i f@{e3, e6, e3, e4}} /. f → fff
Total@%
```

```
Out[2713]= {i e3, e1, -e1, -i e3}
```

```
Out[2714]= 0
```

```
In[2715]:= (* Last reduction in eq. 20 *)
```

```
Expand@Total[ $\frac{i}{2} \{\text{\%\\%\\%\\%, \\%\\%\\%, \\%, \%}\}$ ]
```

```
Out[2715]= 2 e3 - 2 i e1
```

```
In[2716]:= (* Last reduction in eq. 21 *)
```

```
i {i f@{e2, e5}, -f@{e5, e6}, -f@{e2, e4}, -i f@{e4, e6}} /. f → fff
Total@%
```

```
Out[2716]= {e3, -i e1, -i e1, e3}
```

```
Out[2717]= 2 e3 - 2 i e1
```

The above reduction did not apply the final $\circ v$. If we do that drives the result to zero.

```
In[2718]:= %  $\circ v$ 
```

```
Out[2718]= 0
```

Now we perform the commutation action in one step ($\Lambda 1[-\nu^*] \odot u1r[2,1]$, notice the use of $-\nu^*$ not ν , which would drive the result to zero).

There is also a factors of 2 that is related to the change in basis from 2014 to 2019. Then compare it against $u1g$ and the result from the manual decomposition above:

```
In[2719]:= (**)
Expand[2 (\Lambda[1, -\nu^*] /. f → fff)  $\odot$  (u1r[2, 1] /. f → fff)]
Expand[2 u1g[2, 1] /. f → fff]
Expand[% == %% == %%%]
```

```
Out[2719]= 2 e3 - 2 i e1
```

```
Out[2720]= 2 e3 - 2 i e1
```

```
Out[2721]= True
```

Now test the reverse action ($\Lambda 1[\nu] \odot u1g[2,1]$, for which we need to use ν (not $-\nu^*$ as above) which would drive this result to zero):

```
In[2722]:= (* u1 G→R *)
(Δ[1, ν] /. f → fff) ⊙ (u1g[[2, 1]] /. f → fff)
u1r[[2, 1]] /. f → fff
Expand[% == %%]
```

Out[2722]= $-e_4 + i e_5$

```
Out[2723]=  $2\left(-\frac{e_4}{4} + \frac{i e_5}{4}\right) - 2\left(\frac{e_4}{4} - \frac{i e_5}{4}\right)$ 
```

Out[2724]= True

Moving on, we now build the machinery to perform these commutations across the SM particles:

```
In[2725]:= (* Function to output the action with input values of Δ_nν and the bioctonion particle *)
doAction[Δg_, particle_] := Expand[(Δg /. f → fff) ⊙ (particle /. f → fff)];
```

```
In[2726]:= (* u1 R→G *)
```

```
In[2727]:= doAction[Δ[1, -ν*], u1r[[2, 1]]]
u1g[[2, 1]] /. f → fff
Expand[% == %%]
```

Out[2727]= $e_3 - i e_1$

```
Out[2728]=  $4\left(\frac{e_3}{4} - \frac{i e_1}{4}\right)$ 
```

Out[2729]= True

```
In[2730]:= (* u1 G→R *)
```

```
In[2731]:= doAction[Δ1[ν], u1g[[2, 1]]]
u1r[[2, 1]] /. f → fff
Expand[% == %%]
```

Out[2731]= $-e_4 + i e_5$

```
Out[2732]=  $2\left(-\frac{e_4}{4} + \frac{i e_5}{4}\right) - 2\left(\frac{e_4}{4} - \frac{i e_5}{4}\right)$ 
```

Out[2733]= True

Function to output the action across SM particles with a string input identifying Δ_n .

I need more information on the flavor and color transformations involved in Δ_n ,
but this helps deduce that by trying various combinations.

```
In[2734]:= (* Define the list of particle flavors *)
flavorList = Flatten@pTypeList[[;; 2]]
```

Out[2734]= {e, ν, d, u}

```
In[2735]:= doSMaction@in_ := Column[Table[Column[Table[Row[Table[
inExp = ToExpression@in;
prt = ToExpression[flavorList[[flv]] <> ToString@gen <> If[flv > 2, rgb[[col]], " "]];
sppt = ToString[
Subsuperscript[flavorList[[flv]], gen, If[flv > 2, rgb[[col]], " "]], TraditionalForm];
ssprt = Style[sppt, Bold, Red];
(* skip color iterations on leptons *)
If[flv < 3 && col > 1, Nothing,
Column[{ssprt, Row@{sppt <> "=" , Row@{Expand@prt[[2, 1]] /. f → fff}, "      "},
Row@{in <> "O" <> sppt <> "=" ,
doAction[inExp, prt[[2, 1]] /. f → fff], "      "}], Center]],
{col, 3}}],
{flv, 4}], Center],
{gen, 3}]];
```

```
In[2736]:= (* u1 R->G *)
doSMaction@"\Lambda[1,-\nu^*]"
```

$$\begin{aligned}
& \mathbf{e}_1 = -3e_3 - ie_1 \\
& \Lambda[1, -\nu^*] \odot \mathbf{e}_1 = -2e_4 - 2ie_5 \\
& \mathbf{v}_1 = 1 + 0.i \\
& \Lambda[1, -\nu^*] \odot \mathbf{v}_1 = 0 \\
& \mathbf{d}_1^r = -e_3 - ie_1 \quad \mathbf{d}_1^g = 0 \quad \mathbf{d}_1^b = 0 \\
& \Lambda[1, -\nu^*] \odot \mathbf{d}_1^r = -e_4 - ie_5 \quad \Lambda[1, -\nu^*] \odot \mathbf{d}_1^g = 0 \quad \Lambda[1, -\nu^*] \odot \mathbf{d}_1^b = 0 \\
& \mathbf{u}_1^r = -e_4 + ie_5 \quad \mathbf{u}_1^g = e_3 - ie_1 \quad \mathbf{u}_1^b = -\frac{e_7}{2} \\
& \Lambda[1, -\nu^*] \odot \mathbf{u}_1^r = e_3 - ie_1 \quad \Lambda[1, -\nu^*] \odot \mathbf{u}_1^g = 0 \quad \Lambda[1, -\nu^*] \odot \mathbf{u}_1^b = \frac{e_7}{4}
\end{aligned}$$

$$\begin{aligned}
& \mathbf{e}_2 = 2e_6 + 2ie_2 \\
& \Lambda[1, -\nu^*] \odot \mathbf{e}_2 = 0
\end{aligned}$$

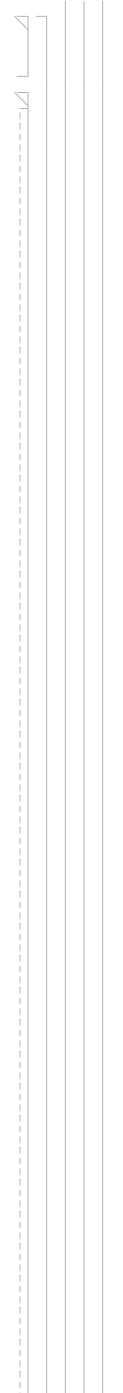
$$\Lambda[1, -\nu^*] \odot \mathbf{v}_2 = \frac{3e_2}{2} - \frac{3ie_6}{2}$$

$$\begin{aligned}
& \mathbf{d}_2^r = 2e_6 + 2ie_2 \quad \mathbf{d}_2^g = 2e_7 + 2i \quad \mathbf{d}_2^b = -2e_5 - 2ie_4 \\
& \Lambda[1, -\nu^*] \odot \mathbf{d}_2^r = 0 \quad \Lambda[1, -\nu^*] \odot \mathbf{d}_2^g = e_2 - ie_6 \quad \Lambda[1, -\nu^*] \odot \mathbf{d}_2^b = 2e_1 + 2ie_3 \\
& \mathbf{u}_2^r = -e_4 + ie_5 \quad \mathbf{u}_2^g = e_3 - ie_1 \quad \mathbf{u}_2^b = -\frac{e_7}{2} \\
& \Lambda[1, -\nu^*] \odot \mathbf{u}_2^r = e_3 - ie_1 \quad \Lambda[1, -\nu^*] \odot \mathbf{u}_2^g = 0 \quad \Lambda[1, -\nu^*] \odot \mathbf{u}_2^b = \frac{e_7}{4}
\end{aligned}$$

$$\begin{aligned}
& \mathbf{e}_3 = 2e_5 + 2ie_4 \\
& \Lambda[1, -\nu^*] \odot \mathbf{e}_3 = -2e_1 - 2ie_3
\end{aligned}$$

$$\Lambda[1, -\nu^*] \odot \mathbf{v}_3 = 0$$

$$\begin{aligned}
& \mathbf{d}_3^r = 0 \quad \mathbf{d}_3^g = 0 \quad \mathbf{d}_3^b = e_3 + ie_1 \\
& \Lambda[1, -\nu^*] \odot \mathbf{d}_3^r = 0 \quad \Lambda[1, -\nu^*] \odot \mathbf{d}_3^g = 0 \quad \Lambda[1, -\nu^*] \odot \mathbf{d}_3^b = e_4 + ie_5 \\
& \mathbf{u}_3^r = 0 \quad \mathbf{u}_3^g = -e_3 - ie_1 \quad \mathbf{u}_3^b = 0 \\
& \Lambda[1, -\nu^*] \odot \mathbf{u}_3^r = 0 \quad \Lambda[1, -\nu^*] \odot \mathbf{u}_3^g = -e_4 - ie_5 \quad \Lambda[1, -\nu^*] \odot \mathbf{u}_3^b = 0
\end{aligned}$$



Visualize the entire family of flavors across 3 generations (2014 version)

These show the unevaluated form as well as the octonion reduced form, with the unevaluated (string form) showing the appending of “ $\circ v$ ” along with the evaluation of that:

```
In[2737]:= doSM14 := Column[Table[Column[Table[Row[Table[
  prt = Expand[ToExpression[flavorList[[flv]] <> ToString@gen <> If[flv > 2, rgb[[col]], " "]];
  sprt = ToString[
    Subsuperscript[flavorList[[flv]], gen, If[flv > 2, rgb[[col]], " "]], TraditionalForm];
  ssprt = Style[sprt, Bold, Red];
  sprt1 = ToString[Expand[prt[[2, 1]]], TraditionalForm] <> "      ";
  sprt2 = ToString[Expand[prt[[2, 1]] /. f → fff], TraditionalForm];
  sprt3 = sprt <> "\circ v" <> ToString[Expand[(prt[[2, 1]] /. f → fff) \circ v], TraditionalForm];
  (* skip color iterations on leptons *)
  If[flv < 3 && col > 1, Nothing,
    Column[{ssprt, sprt1, sprt2, sprt3}, Center]],
  {col, 3}]], {flv, 4}], Center],
{gen, 3}], Center];
```

In[2738]:= doSM14

$$\begin{aligned}
& \text{f}([e_1, e_2, e_6]) + f([e_1, e_4, e_5]) - i f([e_1]) - f([e_3]) \\
& \quad -3 e_3 - i e_1 \\
& e_1 \circ v = -2 e_3 - 2 i e_1 \\
& \quad 1 - 0. i \\
& \quad 1 - 0. i \\
& v_1 \circ v = 0.5 + (0. + 0.5 i) e_7 \\
& \text{d}\bar{e}_1 \\
& f([e_1, e_2, e_6]) - f([e_1, e_4, e_5]) - i f([e_1]) - f([e_3]) \\
& \quad -e_3 - i e_1 \\
& d\bar{e}_1 \circ v = -e_3 - i e_1 \\
& \text{d}\bar{e}_2 \\
& f([e_1, e_3, e_4]) - i f([e_1, e_3, e_5]) + i f([e_4]) + f([e_5]) \\
& \quad -i f([e_1, e_2, e_4]) - f([e_1, e_2, e_5]) - f([e_1, e_4, e_6]) + i f([e_1, e_5, e_6]) \\
& e_3 - i e_1 \\
& d\bar{e}_2 \circ v = 0 \\
& \text{d}\bar{e}_3 \\
& f([e_1, e_2, e_6]) + 2 i f([e_2, e_2, e_6]) - 2 i f([e_2, e_2, e_6]) - 2 f([e_2, e_2, e_6]) \\
& \quad 2 i f([e_2, e_2, e_6]) + 2 f([e_2, e_2, e_6]) - 2 f([e_2, e_2, e_6]) + 2 i f([e_2, e_2, e_6]) \\
& e_3 - i e_1 \\
& u\bar{e}_3 \circ v = 0 \\
& \text{d}\bar{e}_4 \\
& f([e_1, e_2, e_5]) + i f([e_1, e_3, e_5]) + i f([e_2]) + f([e_6]) \\
& \quad 2 e_6 + 2 i e_2 \\
& e_2 \circ v = 2 e_6 + 2 i e_2 \\
& \text{d}\bar{e}_4 \\
& f([e_1, e_2, e_5]) + i f([e_1, e_3, e_5]) + f([e_2]) + f([e_6]) \\
& \quad -i f([e_1, e_2, e_4]) - f([e_1, e_2, e_5]) + f([e_1, e_4, e_6]) - i f([e_1, e_5, e_6]) \\
& \quad 2 e_7 + 2 i \\
& d\bar{e}_4 \circ v = 2 e_6 + 2 i e_2 \\
& \text{d}\bar{e}_5 \\
& f([e_1, e_2, e_6]) + 2 i f([e_2, e_2, e_6]) - 2 i f([e_2, e_2, e_6]) - 2 f([e_2, e_2, e_6]) \\
& \quad 2 i f([e_2, e_2, e_6]) + 2 f([e_2, e_2, e_6]) - 2 f([e_2, e_2, e_6]) + 2 i f([e_2, e_2, e_6]) \\
& e_3 - i e_1 \\
& u\bar{e}_5 \circ v = 0 \\
& \text{d}\bar{e}_6 \\
& -f([e_1, e_3, e_4]) + i f([e_1, e_3, e_5]) + i f([e_4]) + f([e_5]) \\
& \quad 2 e_5 + 2 i e_4 \\
& e_2 \circ v = 2 e_5 + 2 i e_4 \\
& \text{d}\bar{e}_6 \\
& -i f([e_1, e_2, e_4]) - f([e_1, e_2, e_5]) - f([e_1, e_4, e_6]) - i f([e_1, e_5, e_6]) \\
& \quad 2 i \\
& v_2 \circ v = -e_7 + i \\
& \text{d}\bar{e}_7 \\
& -i f([e_1, e_2, e_4]) + f([e_1, e_2, e_5]) + f([e_1, e_4, e_6]) + i f([e_1, e_5, e_6]) \\
& \quad 0 \\
& d\bar{e}_7 \circ v = 0 \\
& \text{d}\bar{e}_8 \\
& f([e_1, e_3, e_4]) + i f([e_1, e_3, e_5]) + f([e_4]) + f([e_5]) \\
& \quad -e_3 - i e_1 \\
& u\bar{e}_8 \circ v = -e_3 - i e_1 \\
& \text{d}\bar{e}_9 \\
& -f([e_1, e_2, e_3]) - i f([e_1, e_3, e_4]) + f([e_2]) + f([e_3]) \\
& \quad -e_3 + i e_1 \\
& u\bar{e}_9 \circ v = 0 \\
& \text{d}\bar{e}_{10} \\
& f([e_1, e_3, e_4]) + i f([e_1, e_3, e_5]) - f([e_1, e_4, e_6]) + f([e_5]) \\
& \quad 0 \\
& u\bar{e}_{10} \circ v = 0
\end{aligned}$$

Visualize the entire family of flavors across 3 generations (2019 version)

These show the reduced octonion structure with pre/post sS, the Q based on the sS used (not sure if this is what is intended), as well as the unevaluated (string form) showing what I assume is the commutation structure of eqs. 19, 22 and 23 action to derive proper charge Q (which doesn't seem to work at all). I need clarification on how this is supposed to operate (maybe more examples):

```
In[2739]:= doSM := Column[Table[Column[Table[Row[Table[
  prt = Expand[
    ToExpression[flavorList[[flv]] <> ToString@gen <> If[flv > 2, rgb[[col]], " "]] /. f → fff];
  sprt = ToString[Subsuperscript[flavorList[[flv]], gen, If[flv > 2, rgb[[col]], " "]],
    TraditionalForm];
  ssprt = Style[sprt, Bold, Red];
  sQ = su3Q@prt;
  sQ1 = "Q=" <> ToString[sQ[[2]], TraditionalForm];
  sQ2 = ToString[Row@sQ, TraditionalForm] <> "⊖" <> ToString[prt[[1, 2]], TraditionalForm] <>
    sprt <> ToString[prt[[1, 1]], TraditionalForm] <> "=";
  sQ3 = "i(" <> ToString[sQ[[2]], TraditionalForm] <> ToString[Row@prt[[1]], TraditionalForm] <>
    "⊖" <> ToString[Row@prt[[1]], TraditionalForm] <> sprt <>
    ToString[Row@prt[[3]], TraditionalForm] <> ")=";
  sQ4 = ToString[Row@sQ[[;; 2]], TraditionalForm] <> sprt;
  (* skip color iterations on leptons *)
  If[flv < 3 && col > 1, Nothing,
    Column[{ssprt, Row@prt, sQ1, Row@sQ2, sQ3, sQ4, "      "}], Center]],
  {col, 3}]],
  {flv, 4}], Center],
{gen, 3}]], Center];
```

In[2740]:= **doSM**

\mathbf{e}_1 $\{s, S^*\}\{-3 e_3 - i e_1\}\{s^*, S^*\}$ $Q=\frac{2}{3}$ $i \frac{2}{3} s \odot S^* e_1 s = i(\frac{2}{3} s S^* \odot s S^* e_1 s^* S) = i \frac{2}{3} e_1$	\mathbf{v}_1 $\{s^*, S\}\{1. + 0. i\}\{s^*, S\}$ $Q=\frac{1}{3}$ $i \frac{1}{3} s^* \odot S v_1 s^* = i(\frac{1}{3} s^* S \odot s^* S v_1 s^* S) = i \frac{1}{3} v_1$	
\mathbf{d}_1^r $\{s, S^*\}\{-e_3 - i e_1\}\{s^*, S\}$ $Q=\frac{2}{3}$ $i \frac{2}{3} s \odot S^* d_1^r s = i(\frac{2}{3} s S^* \odot s S^* d_1^r s^* S) = i \frac{2}{3} d_1^r$	\mathbf{d}_1^g $\{s, S^*\}\{0\}\{s^*, S\}$ $Q=\frac{2}{3}$ $i \frac{2}{3} s \odot S^* d_1^g s = i(\frac{2}{3} s S^* \odot s S^* d_1^g s^* S) = i \frac{2}{3} d_1^g$	\mathbf{d}_1^b $\{s, S^*\}\{0\}\{s^*, S\}$ $Q=\frac{2}{3}$ $i \frac{2}{3} s \odot S^* d_1^b s = i(\frac{2}{3} s S^* \odot s S^* d_1^b s^* S) = i \frac{2}{3} d_1^b$
\mathbf{u}_1^r $\{s, S^*\}\{-e_4 + i e_5\}\{s, S\}$ $Q=\frac{2}{3}$ $i \frac{2}{3} s \odot S^* u_1^r s = i(\frac{2}{3} s S^* \odot s S^* u_1^r s S) = i \frac{2}{3} u_1^r$	\mathbf{u}_1^g $\{s, S^*\}\{e_3 - i e_1\}\{s, S\}$ $Q=\frac{2}{3}$ $i \frac{2}{3} s \odot S^* u_1^g s = i(\frac{2}{3} s S^* \odot s S^* u_1^g s S) = i \frac{2}{3} u_1^g$	\mathbf{u}_1^b $\{s, S^*\}\{-\frac{e_2}{2}\}\{s, S\}$ $Q=\frac{2}{3}$ $i \frac{2}{3} s \odot S^* u_1^b s = i(\frac{2}{3} s S^* \odot s S^* u_1^b s S) = i \frac{2}{3} u_1^b$
\mathbf{e}_2 $\{s, S^*\}\{2 e_6 + 2 i e_2\}\{s^*, S^*\}$ $Q=\frac{2}{3}$ $i \frac{2}{3} s \odot S^* e_2 s = i(\frac{2}{3} s S^* \odot s S^* e_2 s^* S) = i \frac{2}{3} e_2$		
\mathbf{v}_2 $\{s, S^*\}\{3 e_7\}\{s, S^*\}$ $Q=\frac{2}{3}$ $i \frac{2}{3} s \odot S^* v_2 s = i(\frac{2}{3} s S^* \odot s S^* v_2 s S^*) = i \frac{2}{3} v_2$		
\mathbf{d}_2^r $\{s, S^*\}\{2 e_6 + 2 i e_2\}\{s^*, S\}$ $Q=\frac{2}{3}$ $i \frac{2}{3} s \odot S^* d_2^r s = i(\frac{2}{3} s S^* \odot s S^* d_2^r s^* S) = i \frac{2}{3} d_2^r$	\mathbf{d}_2^g $\{s, S^*\}\{2 e_7 + 2 i\}\{s^*, S\}$ $Q=\frac{2}{3}$ $i \frac{2}{3} s \odot S^* d_2^g s = i(\frac{2}{3} s S^* \odot s S^* d_2^g s^* S) = i \frac{2}{3} d_2^g$	\mathbf{d}_2^b $\{s, S^*\}\{-2 e_5 - 2 i e_4\}\{s^*, S\}$ $Q=\frac{2}{3}$ $i \frac{2}{3} s \odot S^* d_2^b s = i(\frac{2}{3} s S^* \odot s S^* d_2^b s^* S) = i \frac{2}{3} d_2^b$
\mathbf{u}_2^r $\{s^*, S^*\}\{-e_4 + i e_5\}\{s^*, S\}$ $Q=1$ $i 1 s^* \odot S^* u_2^r s^* = i(1 s^* S^* \odot s^* S^* u_2^r s^* S) = i 1 u_2^r$	\mathbf{u}_2^g $\{s^*, S^*\}\{e_3 - i e_1\}\{s^*, S\}$ $Q=1$ $i 1 s^* \odot S^* u_2^g s^* = i(1 s^* S^* \odot s^* S^* u_2^g s^* S) = i 1 u_2^g$	\mathbf{u}_2^b $\{s^*, S^*\}\{-\frac{e_2}{2}\}\{s^*, S\}$ $Q=1$ $i 1 s^* \odot S^* u_2^b s^* = i(1 s^* S^* \odot s^* S^* u_2^b s^* S) = i 1 u_2^b$
\mathbf{e}_3 $\{s, S^*\}\{2 e_5 + 2 i e_4\}\{s^*, S^*\}$ $Q=\frac{2}{3}$ $i \frac{2}{3} s \odot S^* e_3 s = i(\frac{2}{3} s S^* \odot s S^* e_3 s^* S) = i \frac{2}{3} e_3$		
\mathbf{v}_3 $\{s^*, S^*\}\{2 i\}\{s, S\}$ $Q=1$ $i 1 s^* \odot S^* v_3 s^* = i(1 s^* S^* \odot s^* S^* v_3 s S) = i 1 v_3$		
\mathbf{d}_3^r $\{s, S^*\}\{0\}\{s^*, S\}$ $Q=\frac{2}{3}$ $i \frac{2}{3} s \odot S^* d_3^r s = i(\frac{2}{3} s S^* \odot s S^* d_3^r s^* S) = i \frac{2}{3} d_3^r$	\mathbf{d}_3^g $\{s, S^*\}\{0\}\{s^*, S\}$ $Q=\frac{2}{3}$ $i \frac{2}{3} s \odot S^* d_3^g s = i(\frac{2}{3} s S^* \odot s S^* d_3^g s^* S) = i \frac{2}{3} d_3^g$	\mathbf{d}_3^b $\{s, S^*\}\{e_3 + i e_1\}\{s^*, S\}$ $Q=\frac{2}{3}$ $i \frac{2}{3} s \odot S^* d_3^b s = i(\frac{2}{3} s S^* \odot s S^* d_3^b s^* S) = i \frac{2}{3} d_3^b$
\mathbf{u}_3^r $\{s, S\}\{0\}\{s^*, S\}$ $Q=\frac{4}{3}$ $i \frac{4}{3} s \odot S u_3^r s = i(\frac{4}{3} s S \odot s S u_3^r s^* S) = i \frac{4}{3} u_3^r$	\mathbf{u}_3^g $\{s, S\}\{-e_3 - i e_1\}\{s^*, S\}$ $Q=\frac{4}{3}$ $i \frac{4}{3} s \odot S u_3^g s = i(\frac{4}{3} s S \odot s S u_3^g s^* S) = i \frac{4}{3} u_3^g$	\mathbf{u}_3^b $\{s, S\}\{0\}\{s^*, S\}$ $Q=\frac{4}{3}$ $i \frac{4}{3} s \odot S u_3^b s = i(\frac{4}{3} s S \odot s S u_3^b s^* S) = i \frac{4}{3} u_3^b$

Out[2740]=

Checking all SM particle color and flavor changing actions (e.g. $\Lambda_n[v] \odot u_1 g$) using 4 forms of $\Lambda_{n=1-8}$ and $g_{n=9-14}$ (i.e. with $f \rightarrow v$, \bar{v} , v^* , v^\dagger)

The output is extensive and given my open questions on the formalism presented, the accuracy likely deviates from the intent of [5], but it is interesting to show how everything transforms. If no transform is found for a particular action, it outputs an * for that action. If a color or flavor changing transformation action is found, it identifies that action with the list of particles the transformation applies to. Note: it only identifies a transformed particle if the source particle has a non-zero reduced value and the resulting match is exact (red) or a \pm integer factor of that particle (blue).

```
In[2741]:= aglist = {"v", "i v", "v*", "v†"};  
  
In[2742]:= (* This generates a list of particle reductions  
    used to compare color and flavor changes against *)  
listSM = Flatten[Table[Table[Table[Table[  
    prt = ToExpression[flavorList[[flv]] <> ToString@gen <> If[flv > 2, rgb[[col]], " "]];  
    (* Conjugate if anti=True *)  
    If[anti, prt = prt*];  
    sprt = ToString[  
        Subsuperscript[flavorList[[flv]], gen, If[flv > 2, rgb[[col]], " "]], TraditionalForm];  
    If[anti, sprt = sprt];  
    ssprt = Style[sprt, Bold, Red];  
    (* skip color iterations on leptons *)  
    If[flv < 3 && col > 1, Nothing,  
        {sprt, Expand[prt[[2, 1]] /. f → fff]}],  
        {col, 3}],  
        {flv, 4}],  
        {gen, 3}],  
        {anti, {False, True}}], 3]
```

e_1	$-i e_1 - 3 e_3$
v_1	$1. + 0. i$
d_1^r	$-i e_1 - e_3$
d_1^g	0
d_1^b	0
u_1^r	$i e_5 - e_4$
u_1^g	$e_3 - i e_1$
u_1^b	$-\frac{e_2}{2}$
e_2	$2 i e_2 + 2 e_6$
v_2	$3 e_7$
d_2^r	$2 i e_2 + 2 e_6$
d_2^g	$2 e_7 + 2 i$
d_2^b	$-2 i e_4 - 2 e_5$
u_2^r	$i e_5 - e_4$
u_2^g	$e_3 - i e_1$
u_2^b	$-\frac{e_2}{2}$
e_3	$2 i e_4 + 2 e_5$
v_3	$2 i$
d_3^r	0
d_3^g	0
d_3^b	$i e_1 + e_3$
u_3^r	0
u_3^g	$-i e_1 - e_3$
u_3^b	0
\bar{e}_1	$i e_1 - 3 e_3$
\bar{v}_1	$1. + 0. i$
\bar{d}_1^r	$i e_1 - e_3$
\bar{d}_1^g	0
\bar{d}_1^b	0
\bar{u}_1^r	$-e_4 - i e_5$
\bar{u}_1^g	$i e_1 + e_3$
\bar{u}_1^b	$-\frac{e_2}{2}$
\bar{e}_2	$2 e_6 - 2 i e_2$
\bar{v}_2	$3 e_7$
\bar{d}_2^r	$2 e_6 - 2 i e_2$
\bar{d}_2^g	$2 e_7 - 2 i$
\bar{d}_2^b	$2 i e_4 - 2 e_5$
\bar{u}_2^r	$-e_4 - i e_5$
\bar{u}_2^g	$i e_1 + e_3$
\bar{u}_2^b	$-\frac{e_2}{2}$
\bar{e}_3	$2 e_5 - 2 i e_4$
\bar{v}_3	$-2 i$
\bar{d}_3^r	0
\bar{d}_3^g	0
\bar{d}_3^b	$e_3 - i e_1$
\bar{u}_3^r	0
\bar{u}_3^g	$i e_1 - e_3$
\bar{u}_3^b	0

```
In[2743]:= (* This checks one non-zero reduced commutation result (e.g.  $\Delta_n[v] \odot u1g$ ) against all of the non-zero reduced particle values generated in lisSM *)
chkRes@res_ := Module[{res2, pos1p, pos2p, pos3p, pos1m, pos2m, pos3m, pos3pos, pos, pos1, f1},
  (* Generate the position lists for each comparison *)
  res2 = Numerator@# & /@ Simplify[(***)res];
  pos1p = Position[listSM[[All, 2]], Expand[1 res2]];
  pos2p = Position[listSM[[All, 2]], Expand[2 res2]];
  pos3p = Position[listSM[[All, 2]], Expand[3 res2]];
  pos1m = Position[listSM[[All, 2]], Expand[-1 res2]];
  pos2m = Position[listSM[[All, 2]], Expand[-2 res2]];
  pos3m = Position[listSM[[All, 2]], Expand[-3 res2]];
  f1 = {pos1p, pos2p, pos3p, pos1m, pos2m, pos3m};
  pos = DeleteDuplicates@Flatten@f1;
  pos1 = DeleteDuplicates@Flatten@pos1p;
  (* If the input is not zero and there was a position found,
  output that particle - color code for red as exact match and blue is a ±integer factor *)
  If[res != 0 && Length@pos > 0,
    Style[{listSM[[Flatten@pos, 1]], res},
      If[Length@pos1p > 0 && MemberQ[listSM[[pos1, 2]], res], Red, Blue]],
    Style[{"*", res}, Black]]];

```

```
In[2744]:= (* This generates all of the possible combinations of color
and flavor commutations across the SM and lists the checked results *)
doSMactionAll = Module[{prt, sprt, ssprt, pos, res, in},
  Column[Table[Column[Table[Column[Table[Row@Table[
    prt = ToExpression[flavorList[[flv]] <> ToString@gen <> If[flv > 2, rgb[col], " "]];
    (* Conjugate if anti=True *)
    If[anti, prt = prt*];
    sprt = ToString[
      Subsuperscript[flavorList[[flv]], gen, If[flv > 2, rgb[col], " "]], TraditionalForm];
    If[anti, sprt = ToString[sprt, TraditionalForm]];
    ssprt = Style[sprt, Bold, Red];
    (* skip color iterations on leptons *)
    If[flv < 3 && col > 1, Nothing,
      Column[{ssprt,
        Row@{sprt <> "=", Row@{Expand@prt[[2, 1]] /. f → fff}},
        MatrixForm@Join[
          Table[in = "Δ[" <> ToString[#] <> "," <> Aglist[[Ag]] <> "]",
            res = doAction[ToExpression@in, prt[[2, 1]]];
            Row@{in, chkRes@res},
            {Ag, Length@Aglist}] & /@ Range@8,
          Table[in = "g[" <> ToString[#] <> "," <> Aglist[[Ag]] <> "]",
            res = doAction[ToExpression@in, prt[[2, 1]]];
            Row@{in, chkRes@res},
            {Ag, Length@Aglist}] & /@ Range[9, 14]]}, Center]],
      {col, 3}],
    {flv, 4}], Center],
    {gen, 3}], Center],
    {anti, {False, True}}]],
  Center]]]

```

$$\begin{aligned}
& \text{res} = \frac{\epsilon_1}{\epsilon_1 - i\epsilon_1} \\
& \text{Ag1}[[v1, v2], e4 - i\epsilon 5] \quad \text{Ag1}[[v1, v2], e5 + i\epsilon 4] \quad \text{Ag1}[[v1, v2], 2e4 + 2ie5] \quad \text{Ag1}[[v1, v2], -e4 + ie5] \\
& \text{Ag2}[[v1, v2], 0] \quad \text{Ag2}[[v1, v2], -\frac{i\epsilon 3}{2} + \frac{i\epsilon 5}{2}] \quad \text{Ag2}[[v1, v2], 0] \quad \text{Ag2}[[v1, v2], -\frac{i\epsilon 3}{2} + \frac{i\epsilon 5}{2}] \\
& \text{Ag3}[[v1, v2], -\frac{i\epsilon 3}{2} + \frac{i\epsilon 5}{2}] \quad \text{Ag3}[[v1, v2], -\frac{i\epsilon 3}{2} + \frac{i\epsilon 5}{2}] \quad \text{Ag3}[[v1, v2], 0] \quad \text{Ag3}[[v1, v2], -\frac{i\epsilon 3}{2} + \frac{i\epsilon 5}{2}] \\
& \text{Ag4}[[v1, v2], 0] \quad \text{Ag4}[[v1, v2], 0] \quad \text{Ag4}[[v1, v2], 0] \quad \text{Ag4}[[v1, v2], 0] \\
& \text{Ag5}[[v1, v2], \overline{v3}, \overline{v4}], 2\epsilon 7] \quad \text{Ag5}[[v1, v2], 2\epsilon 7] \quad \text{Ag5}[[v1, v2], \overline{v3}, \overline{v4}, \overline{v2}, \overline{v3}, \epsilon 7] \quad \text{Ag5}[[v1, v2], \overline{v3}, \overline{v4}], -2\epsilon 7] \\
& \text{Ag6}[[v1, v2], 0] \quad \text{Ag6}[[v1, v2], 0] \quad \text{Ag6}[[v1, v2], 0] \quad \text{Ag6}[[v1, v2], 0] \\
& \text{Ag7}[[v1, v2], 0] \quad \text{Ag7}[[v1, v2], 0] \quad \text{Ag7}[[v1, v2], 0] \quad \text{Ag7}[[v1, v2], 0] \\
& \text{Ag8}[[v1, v2], -\frac{i\epsilon 3}{2} + \frac{i\epsilon 5}{2}] \quad \text{Ag8}[[v1, v2], -\frac{i\epsilon 3}{2} + \frac{i\epsilon 5}{2}] \quad \text{Ag8}[[v1, v2], 0] \quad \text{Ag8}[[v1, v2], -\frac{i\epsilon 3}{2} + \frac{i\epsilon 5}{2}] \\
& g1[[v1, v2], \overline{v3}, \overline{v4}], \frac{c_4}{36} + \frac{ic_5}{36}] \quad g1[[v1, v2], \overline{v3}, \overline{v4}], -\frac{c_4}{36} - \frac{ic_5}{36}] \quad g1[[v1, v2], \overline{v3}, \overline{v4}], \frac{c_4}{36} - \frac{ic_5}{36}] \quad g1[[v1, v2], \overline{v3}, \overline{v4}], -\frac{c_4}{36} - \frac{ic_5}{36}] \\
& g1[[10, v1], [\overline{v2}, \overline{v3}], \frac{c_4}{36} + \frac{ic_5}{36}] \quad g1[[10, v1], [\overline{v2}, \overline{v3}], -\frac{c_4}{36} - \frac{ic_5}{36}] \quad g1[[10, v1], [\overline{v2}, \overline{v3}], \frac{c_4}{36} - \frac{ic_5}{36}] \quad g1[[10, v1], [\overline{v2}, \overline{v3}], -\frac{c_4}{36} - \frac{ic_5}{36}] \\
& g1[[11, v1], -\frac{ic_7}{72}] \quad g1[[11, v1], \frac{ic_7}{72}] \quad g1[[11, v1], \frac{ic_7}{36}] \quad g1[[11, v1], -\frac{ic_7}{36}] \\
& g1[[12, v1], [v2, v1], \overline{v3}, \overline{v4}], -\frac{c_4}{36}] \quad g1[[12, v1], [v2, v1], \overline{v3}, \overline{v4}], \frac{c_4}{36} - \frac{ic_5}{36}] \quad g1[[12, v1], [v2, v1], \overline{v3}, \overline{v4}], \frac{c_4}{36} + \frac{ic_5}{36}] \quad g1[[12, v1], [v2, v1], \overline{v3}, \overline{v4}], -\frac{c_4}{36}] \\
& g1[[13, v1], [\overline{v2}, \overline{v3}], \frac{c_4}{12} - \frac{ic_5}{12}] \quad g1[[13, v1], [\overline{v2}, \overline{v3}], -\frac{c_4}{12} + \frac{ic_5}{12}] \quad g1[[13, v1], [\overline{v2}, \overline{v3}], \frac{c_4}{6} + \frac{ic_5}{6}] \quad g1[[13, v1], [\overline{v2}, \overline{v3}], -\frac{c_4}{12} + \frac{ic_5}{12}] \\
& g1[[14, v1], [\overline{v2}, \overline{v3}], \frac{c_4}{36} - \frac{ic_5}{36}] \quad g1[[14, v1], [\overline{v2}, \overline{v3}], \frac{c_4}{36} + \frac{ic_5}{36}] \quad g1[[14, v1], [\overline{v2}, \overline{v3}], \frac{c_4}{72} + \frac{ic_5}{72}] \quad g1[[14, v1], [\overline{v2}, \overline{v3}], \frac{c_4}{36} + \frac{ic_5}{36}]
\end{aligned}$$

	$v_3 = 2i$	$v_3 = -2i$
$A1[1,j]v_i, 0$	$A1[1,j]v_i, 0$	$A1[1,j^2]v_i, 0$
$A1[2,j]v_i, 0$	$A1[2,j]v_i, 0$	$A1[2,j^2]v_i, 0$
$A1[3,j]v_i, 0$	$A1[3,j]v_i, 0$	$A1[3,j^2]v_i, 0$
$A1[4,j]v_i, 0$	$A1[4,j]v_i, 0$	$A1[4,j^2]v_i, 0$
$A1[5,j]v_i, 0$	$A1[5,j]v_i, 0$	$A1[5,j^2]v_i, 0$
$A1[6,j]v_i, 0$	$A1[6,j]v_i, 0$	$A1[6,j^2]v_i, 0$
$A1[7,j]v_i, 0$	$A1[7,j]v_i, 0$	$A1[7,j^2]v_i, 0$
$A1[8,j]v_i, 0$	$A1[8,j]v_i, 0$	$A1[8,j^2]v_i, 0$
$g[1,0,j]v_i, 0$	$g[1,0,j]v_i, 0$	$g[1,0,j^2]v_i, 0$
$g[10,0,j]v_i, 0$	$g[10,0,j]v_i, 0$	$g[10,0,j^2]v_i, 0$
$g[11,0,j]v_i, 0$	$g[11,0,j]v_i, 0$	$g[11,0,j^2]v_i, 0$
$g[12,0,j]v_i, 0$	$g[12,0,j]v_i, 0$	$g[12,0,j^2]v_i, 0$
$g[13,0,j]v_i, 0$	$g[13,0,j]v_i, 0$	$g[13,0,j^2]v_i, 0$
$g[14,0,j]v_i, 0$	$g[14,0,j]v_i, 0$	$g[14,0,j^2]v_i, 0$

$(\mathbf{g}(14,v), \mathbf{v}, 0)$	$\mathbf{g}(14,x), \mathbf{x}, 0)$	$\mathbf{g}(14,y), \mathbf{y}, 0)$	$\mathbf{g}(14,z), \mathbf{z}, 0)$	$\mathbf{g}(14,w), \mathbf{w}, 0)$
$\mathbf{g}(13,v), \mathbf{v}, 0)$	$\mathbf{g}(13,x), \mathbf{x}, 0)$	$\mathbf{g}(13,y), \mathbf{y}, 0)$	$\mathbf{g}(13,z), \mathbf{z}, 0)$	$\mathbf{g}(13,w), \mathbf{w}, 0)$
$\left \mathbf{g}(14,v) \right _{\{ \mathbf{c}_2, \mathbf{c}_5 \}} \cdot \frac{\mathbf{c}_5}{72} - \frac{\mathbf{c}_2}{72}$	$\left \mathbf{g}(14,x) \right _{\{ \mathbf{c}_2, \mathbf{c}_5 \}} \cdot \frac{\mathbf{c}_5}{72} + \frac{\mathbf{c}_2}{72}$	$\left \mathbf{g}(14,y) \right _{\{ \mathbf{c}_2, \mathbf{c}_5 \}} \cdot \frac{\mathbf{c}_5}{72} - \frac{\mathbf{c}_2}{72}$	$\left \mathbf{g}(14,z) \right _{\{ \mathbf{c}_2, \mathbf{c}_5 \}} \cdot \frac{\mathbf{c}_5}{72} + \frac{\mathbf{c}_2}{72}$	$\left \mathbf{g}(14,w) \right _{\{ \mathbf{c}_2, \mathbf{c}_5 \}} \cdot \frac{\mathbf{c}_5}{72} - \frac{\mathbf{c}_2}{72}$
$\overline{\mathbf{c}_1}$	$\overline{\mathbf{c}_1}$	$\overline{\mathbf{c}_1}$	$\overline{\mathbf{c}_1}$	$\overline{\mathbf{c}_1}$
$A[1,v], v, 2\mathbf{c}_2 - 2\mathbf{c}_5]$	$A[1,x], v, [\mathbf{c}_2, \mathbf{c}_5], 2\mathbf{c}_5 + 2\mathbf{c}_2]$	$A[1,y], v, [\mathbf{c}_2, \mathbf{c}_5], \mathbf{c}_2 + \mathbf{c}_5]$	$A[1,z], v, [-2\mathbf{c}_2 + 2\mathbf{c}_5]$	$A[1,w], v, [-2\mathbf{c}_2 + 2\mathbf{c}_5]$
$A[2,v], v, 0)$	$A[2,x], v, 0)$	$A[2,y], v, 0)$	$A[2,z], v, 0)$	$A[2,w], v, 0)$
$A[3,v], v, [\mathbf{c}_2, \mathbf{c}_5], 2\mathbf{c}_2 - 2\mathbf{c}_5]$	$A[3,x], v, [\mathbf{c}_2, \mathbf{c}_5], 2\mathbf{c}_2 - 2\mathbf{c}_5]$	$A[3,y], v, [\mathbf{c}_2, \mathbf{c}_5], 2\mathbf{c}_2 - 2\mathbf{c}_5$	$A[3,z], v, [\mathbf{c}_2, \mathbf{c}_5], 2\mathbf{c}_2 - 2\mathbf{c}_5$	$A[3,w], v, [\mathbf{c}_2, \mathbf{c}_5], 2\mathbf{c}_2 - 2\mathbf{c}_5$
$A[4,v], v, 0)$	$A[4,x], v, 0)$	$A[4,y], v, 0)$	$A[4,z], v, 0)$	$A[4,w], v, 0)$
$A[5,v], v, [\mathbf{c}_1, \mathbf{c}_2, \mathbf{c}_5], 2\mathbf{c}_2 - 2\mathbf{c}_5, \mathbf{c}_7]$	$A[5,x], v, [\mathbf{c}_1, \mathbf{c}_2, \mathbf{c}_5], 2\mathbf{c}_2 - 2\mathbf{c}_5, \mathbf{c}_7]$	$A[5,y], v, [\mathbf{c}_1, \mathbf{c}_2, \mathbf{c}_5], 2\mathbf{c}_2 - 2\mathbf{c}_5, \mathbf{c}_7]$	$A[5,z], v, [\mathbf{c}_1, \mathbf{c}_2, \mathbf{c}_5], 2\mathbf{c}_2 - 2\mathbf{c}_5, \mathbf{c}_7]$	$A[5,w], v, [\mathbf{c}_1, \mathbf{c}_2, \mathbf{c}_5], 2\mathbf{c}_2 - 2\mathbf{c}_5, \mathbf{c}_7]$
$A[6,v], v, 0)$	$A[6,x], v, 0)$	$A[6,y], v, 0)$	$A[6,z], v, 0)$	$A[6,w], v, 0)$
$A[7,v], v, [\mathbf{c}_2, \mathbf{c}_5], 2\mathbf{c}_2 - 2\mathbf{c}_5]$	$A[7,x], v, [\mathbf{c}_2, \mathbf{c}_5], 2\mathbf{c}_2 - 2\mathbf{c}_5$	$A[7,y], v, [\mathbf{c}_2, \mathbf{c}_5], 2\mathbf{c}_2 - 2\mathbf{c}_5$	$A[7,z], v, [\mathbf{c}_2, \mathbf{c}_5], 2\mathbf{c}_2 - 2\mathbf{c}_5$	$A[7,w], v, [\mathbf{c}_2, \mathbf{c}_5], 2\mathbf{c}_2 - 2\mathbf{c}_5$
$A[8,v], v, [\mathbf{c}_2, \mathbf{c}_5], \frac{\mathbf{c}_5}{2}$	$A[8,x], v, [\mathbf{c}_2, \mathbf{c}_5], \frac{3}{2}\mathbf{c}_2 - \frac{1}{2}\mathbf{c}_5$	$A[8,y], v, [\mathbf{c}_2, \mathbf{c}_5], \frac{1}{2}\mathbf{c}_2 - \frac{3}{2}\mathbf{c}_5$	$A[8,z], v, [\mathbf{c}_2, \mathbf{c}_5], \frac{1}{2}\mathbf{c}_2 - \frac{3}{2}\mathbf{c}_5$	$A[8,w], v, [\mathbf{c}_2, \mathbf{c}_5], \frac{1}{2}\mathbf{c}_2 - \frac{3}{2}\mathbf{c}_5$
$g(v)^2 \left [\mathbf{c}_1, \mathbf{c}_2, \mathbf{c}_5], \frac{\mathbf{c}_5}{2} \right. - \frac{\mathbf{c}_2}{2}$	$g(x)^2 \left [\mathbf{c}_1, \mathbf{c}_2, \mathbf{c}_5], -\frac{\mathbf{c}_5}{2} \right. - \frac{\mathbf{c}_2}{2}$	$g(y)^2 \left [\mathbf{c}_1, \mathbf{c}_2, \mathbf{c}_5], \frac{\mathbf{c}_5}{2} \right. - \frac{\mathbf{c}_2}{2}$	$g(z)^2 \left [\mathbf{c}_1, \mathbf{c}_2, \mathbf{c}_5], -\frac{\mathbf{c}_5}{2} \right. - \frac{\mathbf{c}_2}{2}$	$g(w)^2 \left [\mathbf{c}_1, \mathbf{c}_2, \mathbf{c}_5], -\frac{\mathbf{c}_5}{2} \right. - \frac{\mathbf{c}_2}{2}$
\cdots	\cdots	\cdots	\cdots	\cdots
\cdots	\cdots	\cdots	\cdots	\cdots

$$\begin{aligned}
& g(10,v)[(v_1, v_2, v_3), \frac{v_1}{72}, \frac{v_2}{72}] = g(10,v)[(v_1, v_2, v_3), \frac{v_1}{36}, \frac{v_2}{36}], \quad g(10,v)[(v_1, v_2, v_3), \frac{v_1}{36}, \frac{v_2}{36}] = g(10,v)[(v_1, v_2, v_3), \frac{v_1}{72}, \frac{v_2}{72}], \\
& g(11,v)[(v_1, v_2, v_3), \frac{v_1}{36}] = g(11,v)[(v_1, v_2, v_3), \frac{v_1}{36}], \quad g(11,v)[(v_1, v_2, v_3), \frac{v_1}{36}] = g(11,v)[(v_1, v_2, v_3), \frac{v_1}{36}], \\
& g(11,v)[(v_1, v_2, v_3), \frac{v_1}{36}, \frac{v_2}{36}] = g(11,v)[(v_1, v_2, v_3), \frac{v_1}{36}, \frac{v_2}{36}], \quad g(11,v)[(v_1, v_2, v_3), \frac{v_1}{36}, \frac{v_2}{36}] = g(11,v)[(v_1, v_2, v_3), \frac{v_1}{36}, \frac{v_2}{36}], \\
& g(12,v)[(v_1, v_2, v_3), \frac{v_1}{36}] = g(12,v)[(v_1, v_2, v_3), \frac{v_1}{36}], \quad g(12,v)[(v_1, v_2, v_3), \frac{v_1}{36}] = g(12,v)[(v_1, v_2, v_3), \frac{v_1}{36}], \\
& g(12,v)[(v_1, v_2, v_3), \frac{v_1}{36}, \frac{v_2}{36}] = g(12,v)[(v_1, v_2, v_3), \frac{v_1}{36}, \frac{v_2}{36}], \quad g(12,v)[(v_1, v_2, v_3), \frac{v_1}{36}, \frac{v_2}{36}] = g(12,v)[(v_1, v_2, v_3), \frac{v_1}{36}, \frac{v_2}{36}], \\
& g(13,v)[(v_1, v_2, v_3), \frac{v_1}{6}, \frac{v_2}{6}] = g(13,v)[(v_1, v_2, v_3), \frac{v_1}{12}, \frac{v_2}{12}], \quad g(13,v)[(v_1, v_2, v_3), \frac{v_1}{12}, \frac{v_2}{12}] = g(13,v)[(v_1, v_2, v_3), \frac{v_1}{6}, \frac{v_2}{6}], \\
& g(13,v)[(v_1, v_2, v_3), \frac{v_1}{6}, \frac{v_2}{6}, \frac{v_3}{6}] = g(13,v)[(v_1, v_2, v_3), \frac{v_1}{12}, \frac{v_2}{12}, \frac{v_3}{6}], \quad g(13,v)[(v_1, v_2, v_3), \frac{v_1}{12}, \frac{v_2}{12}, \frac{v_3}{6}] = g(13,v)[(v_1, v_2, v_3), \frac{v_1}{6}, \frac{v_2}{6}, \frac{v_3}{6}], \\
& g(14,v)[(v_1, v_2, v_3), \frac{v_1}{72}, \frac{v_2}{72}] = g(14,v)[(v_1, v_2, v_3), \frac{v_1}{36}, \frac{v_2}{36}], \quad g(14,v)[(v_1, v_2, v_3), \frac{v_1}{36}, \frac{v_2}{36}] = g(14,v)[(v_1, v_2, v_3), \frac{v_1}{72}, \frac{v_2}{72}], \\
& g(14,v)[(v_1, v_2, v_3), \frac{v_1}{72}, \frac{v_2}{72}, \frac{v_3}{6}] = g(14,v)[(v_1, v_2, v_3), \frac{v_1}{36}, \frac{v_2}{36}, \frac{v_3}{6}], \quad g(14,v)[(v_1, v_2, v_3), \frac{v_1}{36}, \frac{v_2}{36}, \frac{v_3}{6}] = g(14,v)[(v_1, v_2, v_3), \frac{v_1}{72}, \frac{v_2}{72}, \frac{v_3}{6}]
\end{aligned}$$

$\frac{\partial \tilde{g}}{\partial t} = -2\pi t + 2i$	$\frac{\partial \tilde{g}}{\partial t} = -2\pi t + 2i$	$\frac{\partial \tilde{g}}{\partial t} = -2\pi t + 2i$	$\frac{\partial \tilde{g}}{\partial t} = -2\pi t + 2i$	$\frac{\partial \tilde{g}}{\partial t} = -2\pi t + 2i$	$\frac{\partial \tilde{g}}{\partial t} = -2\pi t + 2i$
$A[1, 0][\overline{w_1}, -e_2 - i e_6]$	$A[1, 4][\overline{w_1}, \overline{w_2}]$, $e_6 - i e_2$	$A[1, 7][\overline{w_1}, -e_2 + i e_6]$	$A[1, 7][\overline{w_1}, e_2 + i e_6]$	$A[1, 9][\overline{w_1}, -e_2 + 2i e_3]$	$A[1, 9][\overline{w_1}, -e_2 - 2i e_3]$
$A[2, 0][\cdot, 0]$	$A[2, 0][\overline{w_1}, 0]$	$A[2, 2][\overline{w_1}, 0]$	$A[2, 3][\overline{w_1}, 0]$	$A[2, 6][\overline{w_1}, 0]$	$A[2, 7][\overline{w_1}, 0]$
$A[3, 0][\cdot, 0]$	$A[3, 3][\overline{w_1}, 0]$	$A[3, 4][\overline{w_1}, 0]$	$A[3, 5][\overline{w_1}, 0]$	$A[3, 6][\overline{w_1}, -e_4 - i e_5]$	$A[3, 7][\overline{w_1}, \overline{w_2}], e_4 + i e_5$
$A[4, 0][\cdot, 0]$	$A[4, 4][\overline{w_1}, 0]$	$A[4, 6][\overline{w_1}, 0]$	$A[4, 7][\overline{w_1}, 0]$	$A[4, 8][\overline{w_1}, 0]$	$A[4, 9][\overline{w_1}, 0]$
$A[5, 3][\overline{w_1}, \overline{w_2}, \overline{w_3}, \overline{w_4}], e_2 - e_1$	$A[5, 4][\overline{w_1}, \overline{w_2}, \overline{w_3}, \overline{w_4}], e_1 + e_2$	$A[5, 5][\overline{w_1}, \overline{w_2}, \overline{w_3}, \overline{w_4}], e_2 + e_1$	$A[5, 7][\overline{w_1}, \overline{w_2}, \overline{w_3}, \overline{w_4}], -e_2 + e_1$	$A[5, 8][\overline{w_1}, \overline{w_2}, \overline{w_3}, \overline{w_4}], -e_4 - i e_5$	$A[5, 9][\overline{w_1}, \overline{w_2}, \overline{w_3}, \overline{w_4}], e_4 + i e_5$
$A[6, 0][\cdot, 0]$	$A[6, 6][\overline{w_1}, 0]$	$A[6, 7][\overline{w_1}, 0]$	$A[6, 8][\overline{w_1}, 0]$	$A[6, 9][\overline{w_1}, 0]$	$A[6, 10][\overline{w_1}, 0]$
$A[7, 0][\cdot, 0]$	$A[7, 4][\overline{w_1}, 0]$	$A[7, 5][\overline{w_1}, 0]$	$A[7, 6][\overline{w_1}, 0]$	$A[7, 7][\overline{w_1}, 0]$	$A[7, 8][\overline{w_1}, 0]$
$A[8, 0][\cdot, 0]$	$A[8, 2][\overline{w_1}, 0]$	$A[8, 3][\overline{w_1}, 0]$	$A[8, 4][\overline{w_1}, 0]$	$A[8, 5][\overline{w_1}, -e_4 - i e_5]$	$A[8, 6][\overline{w_1}, \overline{w_2}], e_4 - i e_5$
$g[9^{(1)}][\overline{w_1}, -e_2 + i e_6]$	$g[9^{(1)}][\overline{w_1}, -e_2 - 2i e_3]$	$g[9^{(1)}][\overline{w_1}, -e_2 + 2i e_3]$	$g[9^{(1)}][\overline{w_1}, -e_2 - 2i e_3]$	$g[9^{(1)}][\overline{w_1}, -e_2 + 2i e_3]$	$g[9^{(1)}][\overline{w_1}, -e_2 - 2i e_3]$
$g[9^{(2)}][\overline{w_1}, \overline{w_2}, -e_2 - 2i e_3]$	$g[9^{(2)}][\overline{w_1}, \overline{w_2}, -e_2 + 2i e_3]$	$g[9^{(2)}][\overline{w_1}, \overline{w_2}, e_2 - 2i e_3]$	$g[9^{(2)}][\overline{w_1}, \overline{w_2}, e_2 + 2i e_3]$	$g[9^{(2)}][\overline{w_1}, \overline{w_2}, -e_2 + 2i e_3]$	$g[9^{(2)}][\overline{w_1}, \overline{w_2}, e_2 - 2i e_3]$
$g[10^{(1)}][\overline{w_1}, -e_2 - \frac{e_4}{2} - \frac{e_2}{2}]$	$g[10^{(1)}][\overline{w_1}, \overline{w_2}, -e_2 - \frac{e_4}{2} + \frac{e_2}{2}]$	$g[10^{(1)}][\overline{w_1}, \overline{w_2}, e_2 - \frac{e_4}{2} + \frac{e_2}{2}]$	$g[10^{(1)}][\overline{w_1}, \overline{w_2}, e_2 + \frac{e_4}{2} - \frac{e_2}{2}]$	$g[10^{(1)}][\overline{w_1}, \overline{w_2}, -e_2 + \frac{e_4}{2} - \frac{e_2}{2}]$	$g[10^{(1)}][\overline{w_1}, \overline{w_2}, -e_2 - \frac{e_4}{2} - \frac{e_2}{2}]$
$g[11^{(1)}][\overline{w_1}, -e_2 - \frac{e_4}{2} + \frac{e_2}{2}]$	$g[11^{(1)}][\overline{w_1}, \overline{w_2}, -e_2 - \frac{e_4}{2} + \frac{e_2}{2}]$	$g[11^{(1)}][\overline{w_1}, \overline{w_2}, e_2 - \frac{e_4}{2} + \frac{e_2}{2}]$	$g[11^{(1)}][\overline{w_1}, \overline{w_2}, e_2 + \frac{e_4}{2} - \frac{e_2}{2}]$	$g[11^{(1)}][\overline{w_1}, \overline{w_2}, -e_2 + \frac{e_4}{2} - \frac{e_2}{2}]$	$g[11^{(1)}][\overline{w_1}, \overline{w_2}, -e_2 - \frac{e_4}{2} + \frac{e_2}{2}]$
$g[12^{(1)}][\overline{w_1}, -e_2 - \frac{e_4}{2} - \frac{e_2}{2}]$	$g[12^{(1)}][\overline{w_1}, \overline{w_2}, -e_2 - \frac{e_4}{2} - \frac{e_2}{2}]$	$g[12^{(1)}][\overline{w_1}, \overline{w_2}, e_2 - \frac{e_4}{2} - \frac{e_2}{2}]$	$g[12^{(1)}][\overline{w_1}, \overline{w_2}, e_2 + \frac{e_4}{2} - \frac{e_2}{2}]$	$g[12^{(1)}][\overline{w_1}, \overline{w_2}, -e_2 + \frac{e_4}{2} - \frac{e_2}{2}]$	$g[12^{(1)}][\overline{w_1}, \overline{w_2}, -e_2 - \frac{e_4}{2} - \frac{e_2}{2}]$
$g[13^{(1)}][\overline{w_1}, -e_2 - \frac{e_4}{2} + \frac{e_2}{2}]$	$g[13^{(1)}][\overline{w_1}, \overline{w_2}, -e_2 - \frac{e_4}{2} + \frac{e_2}{2}]$	$g[13^{(1)}][\overline{w_1}, \overline{w_2}, e_2 - \frac{e_4}{2} + \frac{e_2}{2}]$	$g[13^{(1)}][\overline{w_1}, \overline{w_2}, e_2 + \frac{e_4}{2} - \frac{e_2}{2}]$	$g[13^{(1)}][\overline{w_1}, \overline{w_2}, -e_2 + \frac{e_4}{2} - \frac{e_2}{2}]$	$g[13^{(1)}][\overline{w_1}, \overline{w_2}, -e_2 - \frac{e_4}{2} + \frac{e_2}{2}]$

$\overline{w_2} = 2 \left(\frac{-c_1}{4} + \frac{ic_2}{4} \right)$	$\overline{w_3} = 2 \left(\frac{c_1}{4} + \frac{ic_2}{4} \right)$	$\overline{w_4} = 2 \left(\frac{c_1}{4} - \frac{ic_2}{4} \right)$	$\overline{w_5} = 2 \left(\frac{-c_1}{4} - \frac{ic_2}{4} \right)$
$\text{Alt}_1.v^{\#}[[\overline{w_1}, \overline{w_2}, \overline{w_3}, \overline{w_4}, \overline{w_5}], -c_1 - ic_1]$	$\text{Alt}_1.v^{\#}[[c_1, -c_1], -c_2]$	$\text{Alt}_1.v^{\#}[[c_1, 0], 0]$	$\text{Alt}_1.v^{\#}[[\overline{w_1}, \overline{w_2}, \overline{w_3}, \overline{w_4}, \overline{w_5}], c_1 + ic_1]$
$\text{Alt}_2.v^{\#}[0]$	$\text{Alt}_2.v^{\#}[0]$	$\text{Alt}_2.v^{\#}[0]$	$\text{Alt}_2.v^{\#}[0]$
$\text{Alt}_3.v^{\#}[[\overline{w_1}, \overline{w_2}, \overline{w_3}, \overline{w_4}, \overline{w_5}], \frac{-c_1}{2} + \frac{ic_2}{2}]$	$\text{Alt}_3.v^{\#}[[\overline{w_1}, \overline{w_2}, \overline{w_3}, \overline{w_4}, \overline{w_5}], \frac{c_1}{2} + \frac{ic_2}{2}]$	$\text{Alt}_3.v^{\#}[[\overline{w_1}, \overline{w_2}, \overline{w_3}, \overline{w_4}, \overline{w_5}], \frac{-c_1}{2} - \frac{ic_2}{2}]$	$\text{Alt}_3.v^{\#}[[\overline{w_1}, \overline{w_2}, \overline{w_3}, \overline{w_4}, \overline{w_5}], \frac{c_1}{2} - \frac{ic_2}{2}]$
$\text{Alt}_4.v^{\#}[0]$	$\text{Alt}_4.v^{\#}[0]$	$\text{Alt}_4.v^{\#}[0]$	$\text{Alt}_4.v^{\#}[0]$
$\text{Alt}_5.v^{\#}[-c_2 + ic_2]$	$\text{Alt}_5.v^{\#}[-c_2]$	$\text{Alt}_5.v^{\#}[-c_2]$	$\text{Alt}_5.v^{\#}[-c_2 + ic_2]$
$\text{Alt}_6.v^{\#}[0]$	$\text{Alt}_6.v^{\#}[0]$	$\text{Alt}_6.v^{\#}[0]$	$\text{Alt}_6.v^{\#}[0]$
$\text{Alt}_7.v^{\#}[0]$	$\text{Alt}_7.v^{\#}[0]$	$\text{Alt}_7.v^{\#}[0]$	$\text{Alt}_7.v^{\#}[0]$
$\text{Alt}_8.v^{\#}[[\overline{w_1}, \overline{w_2}, \overline{w_3}, \overline{w_4}, \overline{w_5}], \frac{-c_1}{2} + \frac{ic_2}{2}]$	$\text{Alt}_8.v^{\#}[[\overline{w_1}, \overline{w_2}, \overline{w_3}, \overline{w_4}, \overline{w_5}], \frac{c_1}{2} + \frac{ic_2}{2}]$	$\text{Alt}_8.v^{\#}[[\overline{w_1}, \overline{w_2}, \overline{w_3}, \overline{w_4}, \overline{w_5}], \frac{-c_1}{2} - \frac{ic_2}{2}]$	$\text{Alt}_8.v^{\#}[[\overline{w_1}, \overline{w_2}, \overline{w_3}, \overline{w_4}, \overline{w_5}], \frac{c_1}{2} - \frac{ic_2}{2}]$
$g(\overline{w_1}).v^{\#}[0]$	$g(\overline{w_2}).v^{\#}[0]$	$g(\overline{w_3}).v^{\#}[0]$	$g(\overline{w_4}).v^{\#}[0]$
$g(\overline{10}).v^{\#}[0]$	$g(\overline{11}).v^{\#}[0]$	$g(\overline{12}).v^{\#}[0]$	$g(\overline{13}).v^{\#}[0]$
$g(\overline{11}).v^{\#}[0]$	$g(\overline{12}).v^{\#}[0]$	$g(\overline{13}).v^{\#}[0]$	$g(\overline{14}).v^{\#}[0]$
$g(\overline{12}).v^{\#}[0]$	$g(\overline{13}).v^{\#}[0]$	$g(\overline{14}).v^{\#}[0]$	$g(\overline{15}).v^{\#}[0]$
$g(\overline{13}).v^{\#}[0]$	$g(\overline{14}).v^{\#}[0]$	$g(\overline{15}).v^{\#}[0]$	$g(\overline{16}).v^{\#}[0]$
$g(\overline{14}).v^{\#}[0]$	$g(\overline{15}).v^{\#}[0]$	$g(\overline{16}).v^{\#}[0]$	$g(\overline{17}).v^{\#}[0]$

$\frac{\partial}{\partial x^1} - \frac{c_1}{2}$		
$A(x^1)^T [v_2, v_1, \underline{v_2}, \underline{v_2}, \underline{v_2}, \underline{v_2}] \cdot \frac{c_1}{4}$	$A(x^1)^T [v_2, v_1, \underline{v_2}, \underline{v_2}, \underline{v_2}, \underline{v_2}] \cdot \frac{c_1}{4}$	
$A(x^2)^T [v_1, 0] \cdot 0$	$A(x^2)^T [v_1, 0] \cdot 0$	
$A(x^3)^T [v_1, \underline{v_2}, \underline{v_2}, \underline{v_2}, \underline{v_2}, \underline{v_2}] \cdot \frac{c_1}{4}$	$A(x^3)^T [v_1, \underline{v_2}, \underline{v_2}, \underline{v_2}, \underline{v_2}, \underline{v_2}] \cdot \frac{c_1}{4}$	
$A(x^4)^T [v_1, 0] \cdot 0$	$A(x^4)^T [v_1, 0] \cdot 0$	
$A(x^5)^T [\underline{v_1}, \underline{v_2}, \underline{v_2}, \underline{v_2}, \underline{v_2}, \underline{v_2}] \cdot \frac{c_1 + c_2}{4}$	$A(x^5)^T [\underline{v_1}, \underline{v_2}, \underline{v_2}, \underline{v_2}, \underline{v_2}, \underline{v_2}] \cdot \frac{c_1 + c_2}{4}$	
$A(x^6)^T [v_1, 0] \cdot 0$	$A(x^6)^T [v_1, 0] \cdot 0$	
$A(x^7)^T [v_1, 0] \cdot 0$	$A(x^7)^T [v_1, 0] \cdot 0$	
$A(x^8)^T [v_2, v_1, \underline{v_2}, \underline{v_2}, \underline{v_2}, \underline{v_2}] \cdot \frac{c_1}{4}$	$A(x^8)^T [v_2, v_1, \underline{v_2}, \underline{v_2}, \underline{v_2}, \underline{v_2}] \cdot \frac{c_1}{4}$	
$g(\mathfrak{g}_1)^T [v_2, v_1, \underline{v_2}, \underline{v_2}, \underline{v_2}, \underline{v_2}] \cdot \frac{c_1}{288}$	$g(\mathfrak{g}_1)^T [v_2, v_1, \underline{v_2}, \underline{v_2}, \underline{v_2}, \underline{v_2}] \cdot \frac{c_1}{288}$	
$g(\mathfrak{g}_2)^T [v_1, \underline{v_2}] \cdot \frac{c_1}{288}$	$g(\mathfrak{g}_2)^T [v_1, \underline{v_2}] \cdot \frac{c_1}{288}$	
$g(\mathfrak{g}_11)^T [\underline{v_1}, \underline{v_2}] \cdot \frac{c_1}{288}$	$g(\mathfrak{g}_{11})^T [\underline{v_1}, \underline{v_2}] \cdot \frac{c_1 + c_2}{288}$	
$g(\mathfrak{g}_12)^T [\underline{v_1}, \underline{v_2}] \cdot \frac{c_1 + c_2}{288}$	$g(\mathfrak{g}_{12})^T [\underline{v_1}, \underline{v_2}] \cdot \frac{c_1 + c_2}{288}$	
$g(\mathfrak{g}_{13})^T [\underline{v_1}, \underline{v_2}, \underline{v_2}, \underline{v_2}, \underline{v_2}, \underline{v_2}] \cdot \frac{c_1 + c_2}{48}$	$g(\mathfrak{g}_{13})^T [\underline{v_1}, \underline{v_2}, \underline{v_2}, \underline{v_2}, \underline{v_2}, \underline{v_2}] \cdot \frac{c_1 + c_2}{48}$	
$g(\mathfrak{g}_14)^T [\underline{v_1}, \underline{v_2}, \underline{v_2}, \underline{v_2}, \underline{v_2}, \underline{v_2}] \cdot \frac{c_1}{48}$	$g(\mathfrak{g}_{14})^T [\underline{v_1}, \underline{v_2}, \underline{v_2}, \underline{v_2}, \underline{v_2}, \underline{v_2}] \cdot \frac{c_1}{48}$	

$A[1..x] \leftarrow s, 2e_1 - 2xe_3$ $A[2..y] \leftarrow 0$ $A[3..x] \leftarrow [a_1 \cdot \vec{w}_1], e_4 + ie_5$ $A[4..y] \leftarrow 0$ $A[5..x] \leftarrow [a_2 \cdot \vec{w}_2, 2e_1 - 2xe_2]$ $A[6..y] \leftarrow 0$ $A[7..x] \leftarrow 0$ $A[8..x] \leftarrow [a_3 \cdot \vec{w}_3], e_4 + ie_5$ $g(9..x) \leftarrow 0$	$A[1..x] \leftarrow s, 2e_3 + 2xe_1$ $A[2..y] \leftarrow 0$ $A[3..x] \leftarrow [a_3 \cdot \vec{w}_3], -e_5 + ie_4$ $A[4..y] \leftarrow 0$ $A[5..x] \leftarrow [a_2 \cdot \vec{w}_2, 2e_2 - 2xe_3]$ $A[6..y] \leftarrow 0$ $A[7..x] \leftarrow 0$ $A[8..x] \leftarrow [a_1 \cdot \vec{w}_1], -e_5 + ie_4$ $g(9..x) \leftarrow 0$
$g(10..x) \leftarrow 0$ $g(11..x) \leftarrow 0$ $g(12..x) \leftarrow 0$ $g(13..x) \leftarrow 0$ $g(14..x) \leftarrow [v_2, v_1, \vec{w}_1^T, \vec{w}_2^T, \frac{v_1}{36}, \frac{v_2}{36}]$	$g(10..x) \leftarrow 0$ $g(11..x) \leftarrow 0$ $g(12..x) \leftarrow 0$ $g(13..x) \leftarrow 0$ $g(14..x) \leftarrow [v_3, v_2, \vec{w}_2^T, \vec{w}_3^T, \frac{v_2}{36}, \frac{v_3}{36}]$

```

A11,v[v], 0 A114,v[v], 0
A12,v[v], 0 A124,v[v], 0
A13,v[v], 0 A134,v[v], 0
A15,v[v], 0 A154,v[v], 0
A16,v[v], 0 A164,v[v], 0
A17,v[v], 0 A174,v[v], 0
A18,v[v], 0 A184,v[v], 0
g11,v[v], 0 g110,v[v], 0
g111,v[v], 0 g1104,v[v], 0
g112,v[v], 0 g1114,v[v], 0
g113,v[v], 0 g1124,v[v], 0
g114,v[v], 0 g1144,v[v], 0

```

$\overline{S_2} = 0$	$\overline{S_3} = 0$	\overline{AII}
$AII[1]*[v_1, 0]$	$AII[1]*[v_0, 0]$	$AII[1]*^2[v_1, 0]$
$AII[2]*[v_0, 0]$	$AII[2]*[v_1, 0]$	$AII[2]*^2[v_1, 0]$
$AII[3]*[v_0, 0]$	$AII[3]*[v_1, 0]$	$AII[3]*^2[v_1, 0]$
$AII[4]*[v_0, 0]$	$AII[4]*[v_1, 0]$	$AII[4]*^2[v_1, 0]$
$AII[5]*[v_0, 0]$	$AII[5]*[v_1, 0]$	$AII[5]*^2[v_1, 0]$
$AII[6]*[v_0, 0]$	$AII[6]*[v_1, 0]$	$AII[6]*^2[v_1, 0]$
$AII[7]*[v_0, 0]$	$AII[7]*[v_1, 0]$	$AII[7]*^2[v_1, 0]$
$AII[8]*[v_0, 0]$	$AII[8]*[v_1, 0]$	$AII[8]*^2[v_1, 0]$
$gI[9]^2, [v_0, 0]$	$gI[9]^2, [v_1, 0]$	$gI[9]^2, [v_1, 0]$
$gI[10], [v_0, 0]$	$gI[10], [v_1, 0]$	$gI[10], [v_1, 0]$
$gI[11], [v_0, 0]$	$gI[11], [v_1, 0]$	$gI[11], [v_1, 0]$
$gI[12], [v_0, 0]$	$gI[12], [v_1, 0]$	$gI[12], [v_1, 0]$

$\text{AU1}^{\text{c}}[v_1, 0]$	$\text{AU1}^{\text{c}}[\overline{v}_1^{\text{c}}, \overline{v}_2^{\text{c}}, \overline{v}_3^{\text{c}}, \dots - \overline{v}_n^{\text{c}}]$
$\text{AU2}^{\text{c}}[v_1, 0]$	$\text{AU2}^{\text{c}}[\overline{v}_1^{\text{c}}, \overline{v}_2^{\text{c}}, \overline{v}_3^{\text{c}}, \dots - \overline{v}_n^{\text{c}}]$
$\text{AU3}^{\text{c}}[v_1, \dots, \frac{v_1}{2}, \dots + \frac{v_3}{2}]$	$\text{AU3}^{\text{c}}[\overline{v}_1^{\text{c}}, \dots, \frac{\overline{v}_1^{\text{c}}}{2}, \dots + \frac{\overline{v}_3^{\text{c}}}{2}]$
$\text{AU4}^{\text{c}}[v_1, 0]$	$\text{AU4}^{\text{c}}[\overline{v}_1^{\text{c}}, 0]$
$5^{\text{c}}\#[(v_1, v_2, v_3, v_4, v_5, -v_7)]$	$\text{AU5}^{\text{c}}[v_1, 0]$
$\text{AU6}^{\text{c}}[v_1, 0]$	$\text{AU6}^{\text{c}}[-v_1, 0]$
$\text{AU7}^{\text{c}}[v_1, 0]$	$\text{AU7}^{\text{c}}[-v_1, 0]$
$\text{AU8}^{\text{c}}[v_1, \dots, \frac{v_1}{2}, \dots + \frac{v_3}{2}]$	$\text{AU8}^{\text{c}}[\overline{v}_1^{\text{c}}, \dots, \frac{\overline{v}_1^{\text{c}}}{2}, \dots + \frac{\overline{v}_3^{\text{c}}}{2}]$
$(\text{AU}^{\text{c}})^{\text{c}}[(v_1, v_2), \frac{v_1}{2}, \dots + \frac{v_3}{2}]$	$\text{g}(\overline{v}_1^{\text{c}})^{\text{c}}[v_1, 0]$
$\text{g}((10, 1)^{\text{c}})[v_1, \overline{v}_2^{\text{c}}], \frac{v_1}{2}, \frac{v_3}{2}$	$\text{g}(10, \overline{v}_1^{\text{c}})^{\text{c}}[v_1, 0]$
$\text{g}(11, 1)^{\text{c}}[v_1, 0]$	$\text{g}(11, \overline{v}_1^{\text{c}})^{\text{c}}[v_1, \frac{v_1 - v_3}{2}]$

IX. Addendum : Towards a complete description

Working with Quaternions

```
In[2745]:= γ1 = I e1; γ2 = I e2; γ3 = I e3; γ4 = I e4; γ5 = I e5; γ6 = I e6; γ7 = I e7; γ8 = J e7; γ9 = K e7;
```

```
In[2746]:= ToQuaternion@γ# & /@ Range@9
```

```
Out[2746]= {Quaternion(0, e1, 0, 0), Quaternion(0, e2, 0, 0), Quaternion(0, e3, 0, 0), Quaternion(0, e4, 0, 0),
Quaternion(0, e5, 0, 0), Quaternion(0, e6, 0, 0), Quaternion(0, e7, 0, 0), Quaternion(0, 0, e7, 0), Quaternion(0, 0, 0, e7)}
```

```
In[2747]:= fff@{γ1, γ2, γ3, γ4, γ5, γ6, γ7}
```

```
Out[2747]= i
```

```
In[2748]:= ToQuaternion[fff@{γ1, γ2, γ3, γ4, γ5, γ6, γ7}]
ToQuaternion@γ8
%% ** %
```

```
Out[2748]= Quaternion(0, 1, 0, 0)
```

```
Out[2749]= Quaternion(0, 0, e7, 0)
```

```
Out[2750]= Quaternion(0, 0, 0, e7)
```

```
In[2751]:= fff@{γ1, γ2, γ3, γ4, γ5, γ6, γ7} ** γ8
```

```
Out[2751]= K e7
```

```
In[2752]:= Table[γn θ γm, {n, 8}, {m, 8}]
```

```
Out[2752]=
```

$$\begin{pmatrix} 2 & 0 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 2 & 0 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 2 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 2 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 2 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 2 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & -2iJ & 0 \\ 0 & 0 & 0 & 0 & 0 & 0 & -2iJ & -2J^2 \end{pmatrix}$$

The Setup for Multi – Actions

```
In[2753]:= P1 = 1/2 (1 - γ9)
```

```
Out[2753]= 1/2 (1 - K e7)
```

```
In[2754]:= P2 = 1/2 (1 + γ9)
```

```
Out[2754]= 1/2 (K e7 + 1)
```

```
In[2755]:= m†[a_, b_] := Sum[ (f@{Pi, a} ** f@{Pi, b} + f@{Pi, b} ** f@{Pi, a†} ), {i, 1, 1} ];
```

```
˜m[a_, b_] := Sum[ (f@{Pi, a} ** f@{Pi, b} + f@{Pi, b} ** f@{Pi, ˜a} ), {i, 1, 1} ];
```

```
In[2757]:= m†[a_, b_] := Sum[ (f@{Pi, a, Pi, b} + f@{Pi, b, Pi, a†} ), {i, 1, 1} ];
```

```
˜m[a_, b_] := Sum[ (f@{Pi, a, Pi, b} + f@{Pi, b, Pi, ˜a} ), {i, 1, 1} ];
```

In[2759]:= $\mathbf{m}\mathbf{t}[\mathbf{v}, \mathbf{b}]$

$$\text{Out}[2759]= f\left(\left\{\frac{1}{2}(1 - K e_7), b, \frac{1}{2}(1 - K e_7), -\frac{1}{2} - \frac{i e_7}{2}\right\}\right) + f\left(\left\{\frac{1}{2}(1 - K e_7), \frac{1}{2}(1 + i e_7), \frac{1}{2}(1 - K e_7), b\right\}\right)$$

In[2760]:= $\mathbf{m}\mathbf{t}[\mathbf{v}, \mathbf{w}]$

$$\text{Out}[2760]= f\left(\left\{\frac{1}{2}(1 - K e_7), \frac{1}{2}(-e_7 + i), \frac{1}{2}(1 - K e_7), -\frac{1}{2} - \frac{i e_7}{2}\right\}\right) + f\left(\left\{\frac{1}{2}(1 - K e_7), \frac{1}{2}(1 + i e_7), \frac{1}{2}(1 - K e_7), \frac{1}{2}(-e_7 + i)\right\}\right)$$

In[2761]:= % /. $\mathbf{f} \rightarrow \mathbf{fff}$

$$\begin{aligned} \text{Out}[2761]= & \frac{1}{16}((-3 + 2i)K^2 e_3 + 3K^2 e_4 - (3 - 2i)K^2 e_5 - (3 - 2i)K^2 e_6 - (2 - i)K^2 e_7 + (3 + i)K^2 + (3K + (3 - 2i))Ke_1 + \\ & (3K + (3 - 2i))Ke_2 + 3Ke_3 + (3 - 2i)Ke_4 + 3Ke_5 + 3Ke_6 + (3 + 3i)Ke_7 + (4 - i)K + 2e_7 - 2i) + \\ & \frac{1}{16}((3 - 2i)K^2 e_3 - 3K^2 e_4 + (3 - 2i)K^2 e_5 + (3 - 2i)K^2 e_6 + (2 - i)K^2 e_7 - (3 + i)K^2 + (-3K - (3 - 2i))Ke_1 + \\ & (-3K - (3 - 2i))Ke_2 - 3Ke_3 - (3 - 2i)Ke_4 - 3Ke_5 - 3Ke_6 - (3 + 3i)Ke_7 - (4 - i)K - 2e_7 + 2i) \end{aligned}$$

In[2762]:= **Simplify**@%

$$\text{Out}[2762]= 0$$

In[2763]:= $\tilde{\mathbf{m}}[\mathbf{v}, \mathbf{w}]$

$$\text{Out}[2763]= f\left(\left\{\frac{1}{2}(1 - K e_7), \frac{1}{2}(-e_7 + i), \frac{1}{2}(1 - K e_7), \frac{1}{2}(\tilde{1} + i e_7)\right\}\right) + f\left(\left\{\frac{1}{2}(1 - K e_7), \frac{1}{2}(1 + i e_7), \frac{1}{2}(1 - K e_7), \frac{1}{2}(-e_7 + i)\right\}\right)$$